

1 **Magnetic signature of large exhumed mantle domains of the Southwest Indian**
2 **Ridge: results from a deep-tow geophysical survey over 0 to 11 Ma old seafloor.**

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9 **ABSTRACT**

10 **We investigate the magnetic signature of ultramafic seafloor in the eastern part of the**
11 **Southwest Indian Ridge (SWIR). There, detachment faulting, continuous over 11 Myr,**
12 **exhumed large areas of mantle derived rocks. These exhumed mantle domains occur in the**
13 **form of a smooth rounded topography with broad ridges locally covered by a thin highly**
14 **discontinuous volcanic carapace. We present high-resolution data combining deep-tow**
15 **magnetics, side-scan sonar images and dredged samples collected within two exhumed mantle**
16 **domains between 62°E and 65°E. We show that, despite an ultra-slow spreading rate, volcanic**
17 **areas within robust magmatic segments are characterized by well defined seafloor spreading**
18 **anomalies. By contrast, the exhumed mantle domains, including a few thin volcanic patches,**
19 **reveal a weak and highly variable magnetic pattern. The analysis of the magnetic properties of**
20 **the dredged samples and careful comparison between the nature of the seafloor, the deep-tow**
21 **magnetic anomalies and the seafloor equivalent magnetization suggest that the serpentized**
22 **peridotites do not carry a sufficiently stable remanent magnetization to produce seafloor**
23 **spreading magnetic anomalies in exhumed mantle domains.**

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Adrien Bronner 6/3/14 11:39
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27	1. INTRODUCTION.....	2	dan 17/3/14 21:00
28	2. GEOLOGICAL BACKGROUND.....		Mis en forme: Anglais (E.U.)
29	3. ACQUISITION AND PROCESSING OF MAGNETIC DATA.....		dan 17/3/14 21:00
30	4. MAGNETIC SIGNAL OVER VOLCANIC SEAFLOOR: A SEAFLOOR SPREADING MODEL.....		Mis en forme: Anglais (E.U.)
31	5. MAGNETIC SIGNAL OVER EXHUMED SERPENTINIZED MANTLE.....	7	dan 17/3/14 21:00
32	5.1 THE WESTERN CORRIDOR.....		Mis en forme: Anglais (E.U.)
33	5.2 THE EASTERN CORRIDOR.....		dan 17/3/14 21:00
34	5.3 MAGNETIC STRUCTURE OF THE DIFFERENT TYPES OF SEAFLOOR.....		Mis en forme: Anglais (E.U.)
35	6. MAGNETIC PROPERTIES OF THE DREDGED SAMPLES.....		dan 17/3/14 21:00
36	7. FORWARD MODELING.....		Mis en forme: Anglais (E.U.)
37	7.1 SEAFLOOR SPREADING MODEL.....		dan 17/3/14 21:00
38	7.2 INDUCED MAGNETIZED MODEL.....		Mis en forme: Anglais (E.U.)
39	8. DEPTH OF THE MAGNETIC SOURCES.....		dan 17/3/14 21:00
40	9. DISCUSSION.....		Mis en forme: Anglais (E.U.)
41	9.1 SEAFLOOR SPREADING ANOMALIES.....		dan 17/3/14 21:00
42	9.2 CONTRIBUTION OF MANTLE DERIVED ROCKS.....		Mis en forme: Anglais (E.U.)
43	9.3 CORRUGATED SEAFLOOR AND THE MAGNETIC SIGNAL.....		dan 17/3/14 21:00
44	9.4 VOLCANIC SEAFLOOR AND THE MAGNETIC SIGNAL.....		Mis en forme: Anglais (E.U.)
45	9.5 MARINE MAGNETIC ANOMALIES AT OCT.....		dan 17/3/14 21:00
46	10. CONCLUSION.....		Mis en forme: Anglais (E.U.)

50 **1. INTRODUCTION**

51 The eastern part of the ultra-slow spreading Southwest Indian Ridge (SWIR) is among the
 52 deepest parts of the oceanic ridge system and represents a melt-poor end-member for this system
 53 (Karson et al., 1987; Cannat et al., 1999; Cannat et al., 2008). In this region, crustal accretion differs
 54 from the conventional seafloor spreading scheme as it occurs at about a 14 mm/a full spreading rate
 55 (Patriat et al., 1997) in the form of magmatic but also non-magmatic processes (Cannat et al., 2006).
 56 In the past two decades, numerous papers have revealed the presence of exhumed mantle-derived
 57 rocks in the oceanic domain (Cannat et al., 1992; Cannat et al., 1995) but mechanisms leading to the
 58 formation of such a peculiar seafloor remain poorly understood. Although it has been proposed that
 59 long-lived detachment faults could often accommodate 50% to 70% (Buck et al., 2005) of the plate
 60 separation over ~3 Myr, the eastern part of the SWIR is currently the only known oceanic area
 61 where continuous mantle exhumation over 11 Myr has been observed (Sauter et al., 2013). There,
 62 detachment faulting associated with no or very little volcanic activity seems to be the only process
 63 producing the oceanic lithosphere. The resulting seafloor, called “smooth seafloor” (Cannat et al.,
 64 2006), is thought to be formed by alternating “flip flop” exhumation faulting (Sauter et al., 2013), a
 65 mechanism that has also been proposed to explain the formation of the “zone of exhumed
 66 continental mantle” (Reston and McDermott, 2011) observed along the ocean-continent transition
 67 (OCT) in the Western Iberia margin.

68 The conventional understanding of seafloor magnetic anomalies is that their source mainly
 69 resides in an upper crustal layer of effusive volcanic rocks (e.g. Harrison, 1987). However, studies at
 70 slow spreading ridges have also suggested a contribution from other lithologies, such as gabbros and
 71 serpentized peridotites (Pariso and Johnson, 1993; Nazarova, 1994; Oufi et al., 2002). A better
 72 understanding of the variability of the amplitude of the magnetic anomalies over exhumed mantle
 73 domains is required to assess the validity of kinematic reconstructions at both ultra-slow-spreading
 74 mid oceanic ridges and magma-poor passive margin systems. In this paper, we investigate the

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80 magnetic signal over large exhumed mantle domains in the easternmost part of the SWIR. We
81 present results from a deep-tow geological-geophysical survey over two areas between 62°E and
82 65°E combining magnetic data, geological mapping from side-scan sonar images (from Sauter et al,
83 2013) and dredge sampling. We examine the magnetic signature over a 0 to 11 Ma old smooth
84 seafloor. The aim is to better understand the complexity of the marine magnetic anomalies observed
85 above the serpentinized mantle rocks exhumed at mid oceanic ridges (Sauter et al., 2008). Finally
86 we discuss the implications of our findings for the understanding of exhumed mantle domains at
87 OCTs of magma-poor rifted margins where the origin and significance of broad zones of chaotic
88 magnetic patterns are discussed (Russell and Whitmarsh, 2003; Sibuet et al., 2007; Bronner et al.,
89 2011; [Tucholke and Sibuet, 2012](#); [Bronner et al., 2012](#)).

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90 2. GEOLOGICAL BACKGROUND

91 A significant change in the Africa-Antarctica relative plate motion occurred between magnetic
92 anomaly C8 and C6 (~24 Ma ago) resulting in a 50% decrease in full spreading rate at the SWIR,
93 from slow (24 mm/a) to ultra-slow (14 mm/a; Patriat et al., 2008). This [ultra-slow](#) spreading rate
94 varies only slightly along the 7700 km ridge axis. By contrast, compilations of geophysical and
95 geochemical data along the SWIR reveal large-scale variations of the density and thermal structure
96 of the axial region (e.g. Cannat et al., 1999; Georgen et al., 2001; Cannat et al., 2008). Unusually
97 cold mantle temperatures and relatively thin crust at the eastern SWIR, in particular, east of the
98 Melville transform fault (60°45'E), are supported by evidence on axis (Cannat et al., 2008) as well
99 as off-axis (Sauter et al., 2011). An eastward decreasing crustal thickness and/or mantle temperature
100 is inferred from gravity data along the SWIR axis (Cannat et al., 1999). It is further supported by
101 geochemical proxies for the degree of partial melting in the mantle (e.g. average of the composition
102 of the sodium content of axial basalts derived from the axial zone) suggesting a progressive
103 eastward decrease of the ridge melt supply (Meyzen et al., 2003; Seyler et al., 2003; Cannat et al.,
104 2008). Thin crust in the easternmost part of the SWIR (3.7 km average crustal thickness) is also
105 confirmed by seismic data (Minshull et al., 2006).

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l10 The easternmost part of the SWIR axial valley displays a ridge segmentation that significantly
l11 differs from what is observed at faster spreading ridges such as the Mid-Atlantic Ridge (MAR).
l12 High-relief ridge segments (>3000 m high) are linked by >100 km long, deep axial sections with
l13 almost no volcanic activity (Sauter et al., 2004). The ridge flanks display the widest known areas of
l14 seafloor with no evidence of a volcanic upper crustal layer (Cannat et al., 2006). This non-volcanic
l15 ocean floor has no equivalent at faster spreading ridges. Cannat et al. (2006) called this seafloor
l16 “smooth seafloor” because it occurs in the form of broad ridges with a smooth, rounded topography
l17 and lacks the telltale hummocky morphologies of submarine volcanism. This non-volcanic seafloor
l18 also lacks the corrugations identified on oceanic core complexes at slow spreading ridges. A few
l19 dredges in the axial valley from earlier cruises suggested that the smooth seafloor is associated with
l20 outcrops of serpentinized mantle-derived peridotites (Cannat et al., 2006). Off-axis dredges and
l21 side-scan sonar imagery confirmed that this smooth seafloor is almost entirely composed of
l22 seawater-altered mantle rocks resulting in serpentinized peridotites that were brought to the surface
l23 by large detachment faults on both sides of the ridge axis (Sauter et al., 2013). The detachment
l24 faults are thought to repeatedly flip polarity and have accommodated nearly 100 % of the plate
l25 divergence for the last 10 Myr (Sauter et al., 2013).

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l26 **3. ACQUISITION AND PROCESSING OF MAGNETIC DATA**

l27 Data presented in this paper were collected during R/V Marion Dufresne cruise MD183 in
l28 October 2010 using a 30 kHz side-scan sonar and a three component (3C) magnetometer carried by
l29 the Towed Ocean Bottom Instrument (TOBI; Flewellen et al., 1993). The survey was divided into
l30 two corridors, a western corridor from 62 to 63°E and an eastern corridor from 64 to 65°E. Seven
l31 profiles were acquired in the western corridor, all above exhumed mantle, and four profiles were
l32 acquired in the eastern corridor, one above volcanic seafloor and three above exhumed mantle rocks
l33 (Fig. 1). TOBI was operated at altitudes of 250-700 m above the seafloor at a tow speed of about 2
l34 knots.

136 The three component magnetic data were corrected for the magnetization of the TOBI vehicle
137 using a scalar calibration procedure (Bronner et al., 2013). The magnetic effect of the vehicle was
138 removed with no recourse to its attitude (pitch, roll or heading) as it is commonly done (Isezaki,
139 1986; Korenaga, 1995), but only using the output of the magnetometer and a model of the scalar
140 intensity of the geomagnetic field (e.g. IGRF). Calibration parameters were thus free from
141 orientation bias (see Bronner et al., 2013) and the estimation of both instrumental miscalibration and
142 removal of the vehicle effect were performed simultaneously. In order to constrain the calibration
143 parameters as much as possible, the geomagnetic field was recorded in a 360° calibration loop in a
144 region of the SWIR where the field was assumed to be constant, and with the most variable possible
145 attitude of the vehicle. Variations of pitch and roll were obtained by successively hauling in and
146 paying out the wire. The magnetic signal of the vehicle was found to be about 3500 nT and reduced
147 to less than 10 nT after calibration. Magnetic data presented in this paper were only corrected via
148 these calibration parameters; no filtering was applied and the quality of the processing was
149 confirmed through a comparison between upward continued data and sea-surface proton
150 magnetometer profiles (Bronner et al., 2013).

151 As magnetic data were acquired along uneven altitudes, we used an equivalent source
152 approach (Dampney, 1969) to invert the magnetic profiles and to perform an upward continuation to
153 a constant observation level. We assume that the measured magnetic anomalies are due to uniformly
154 magnetized dipoles that extend infinitely perpendicular to the spreading and profile direction. The
155 so-called “equivalent layer” is draped on the bathymetry 500 m below the seafloor and
156 magnetization directions are assumed to be parallel to the Earth magnetic field (-60° inclination and
157 -30° declination in this area of the SWIR). Magnetization of the dipoles is then computed in the
158 spatial domain as a single linear inversion to the distances between dipoles and observation points
159 (Bronner et al., 2013). Once the magnetization is obtained, upward continuation is performed by
160 computing the magnetic field due to the equivalent sources at the desired observation level (Fig. 2).
161 Over the volcanic seafloor we assume that a standard homogeneous 500 m layer accounts for the
162 observed magnetic anomalies (e.g. Gee and Kent, 2007). The inferred magnetization values are thus

163 divided by the assumed dipole spacing and layer thickness to yield units of ampere per meter.
164 Magnetizations above exhumed mantle areas are calculated in the same way, although we have little
165 knowledge about the source layer thickness there. These magnetizations have thus to be taken with
166 care and are only presented as a comparison to the volcanic seafloor. Variations of inverted
167 magnetizations over exhumed mantle domains could either result from changes of intrinsic
168 magnetization or from variability in the source thickness.

169 To be consistent, all deep-tow magnetic anomaly profiles displayed in Fig. 2 are upward
170 continued to a constant level of 1200 m below sea level (shallowest depth of the TOBI during the
171 whole survey). 2D magnetic anomaly profiles are represented above seafloor topography in which
172 geological interpretations from side-scan images (from Sauter et al., 2013) are superimposed (Fig.
173 2). As the profiles are about 6 km apart (width of the TOBI side-scan swath) we do not perform 3D
174 inversion or magnetic mapping; instead, we calculate magnetizations along profiles and display
175 them as colored strips of arbitrary width superimposed on the bathymetry (Fig. 3). Identification of
176 magnetic anomalies are based on Sauter et al. (2008).

177 The TOBI 30 kHz side-scan sonar provides 3 m resolution acoustic images of the seafloor.
178 Interpretation of the reflectivity combined with results from dredges leads to the distinction between
179 three types of seafloor (see Sauter et al., 2013): (1) volcanic seafloor, corresponding to highly
180 reflective surfaces composed of volcanic cones (<200 m across) and sinuous scarps characteristic of
181 the presence of pillow lava flows, (2) smooth seafloor, corresponding to smooth and homogeneous
182 topography associated with low and uniform reflectivity and, (3) corrugated seafloor (Cannat et al.,
183 2006) associated with striations comparable to the slip surfaces that are commonly observed at
184 oceanic core complexes of the MAR (Cann et al., 1997). As the sedimentary cover is limited to
185 small patches in this region, the nature of the seafloor below is extrapolated from the surrounding
186 exposed rocks.

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188 4. MAGNETIC SIGNAL OVER VOLCANIC SEAFLOOR: A 189 SEAFLOOR SPREADING MODEL

190 Profile 2-5 was acquired between magnetic anomaly C3A on each flank (see Fig. 1) above
191 exclusively volcanic seafloor associated with relatively thick crust suggested by low Residual
192 Mantle Bouguer Anomalies (RMBA < 20mGal; Cannat et al., 2006). We use it as a reference to
193 calibrate the spreading rate and identify the main polarity reversals. The inverted magnetization
194 values reach around 10 A/m at the axis (resulting in a ~500 nT amplitude for the central anomaly,
195 Fig. 2 and Fig.3b) and 5 A/m off-axis. These values are in agreement with previous observations in
196 this area (Searle and Bralee, 2007) and in another section of the SWIR near 58°E (Hosford et al.,
197 2003). Despite the ultra-slow spreading rate, the main magnetic blocks are well resolved (Fig. 2 and
198 3) and associated with relatively strong magnetic contacts.

199 Seafloor spreading anomalies are modeled using MODMAG (Mendel et al., 2005). A
200 14 mm/a uniform full spreading rate associated with a 500 m thick source layer draped over the
201 topography and a 10 A/m magnetization on-axis decreasing to 5 A/m off-axis enables to reproduce
202 the main magnetic anomalies observed over the volcanic crust at profile 2-5 (Fig. 2). At such ultra-
203 slow spreading rate the identification of the seafloor spreading anomalies is more difficult than for
204 faster spreading ridges because reverse and normal polarity blocks may overlap. Therefore, a 0.7
205 contamination coefficient (Tisseau and Patriat, 1981) is used as a good compromise to both account
206 for contamination between adjacent magnetic blocks with different polarity and preserve the small
207 wavelength anomalies such as anomaly C2 (Fig. 2; profile 2-5). There is a reasonable fit between
208 the observed magnetic anomaly profile and this forward model regarding the central Brunhes
209 anomaly while southern anomaly C2A and northern anomaly C3A are in agreement with previous
210 identifications on sea surface magnetic anomaly profiles (Sauter et al., 2008). Anomaly C2A is not
211 clearly identified on the northern flank. This is consistent with observations from Searle and Bralee
212 (2007) who showed that this polarity reversal was either smaller than predicted or missing in the
213 northern flank in this area. We also suggest that anomaly C2 could account for the two small

214 wavelength events observed on both sides of the central anomaly.

215 5. MAGNETIC SIGNAL OVER EXHUMED SERPENTINIZED 216 MANTLE.

217 5.1 THE WESTERN CORRIDOR

218 The western corridor extends between magnetic anomalies C3A (Cannat et al., 2006), and
219 includes two ~100 km long north-south magnetic profiles 6 km apart ([profiles 1-6 and 1-7](#); Fig. 1)
220 and one short (~30 km) profile that does not cross the axis ([profile 1-5](#); Fig. 1). The magnetic data of
221 the east-west profiles are not presented in this paper because the 2D assumption used for the upward
222 continuation and the computation of the magnetization is unreliable in that case. Therefore, we only
223 use the [side-scan](#) images from these east-west lines to constrain the nature of the seafloor. Careful
224 analysis of bathymetry, [side-scan](#) images and dredge samples suggests that the seafloor in this
225 corridor is exclusively made of wide serpentinitized peridotite ridges topped by thin (<100-200 m
226 thick) volcanic patches (Sauter et al., 2013; Fig. 2 and 3). The axial valley is marked by an
227 unconventional morphology comprising a 2000 m high peridotite ridge, called “Cannibal Ridge”
228 that emerges from the axial domain (Fig. 2).

229 The axial magnetic anomaly is hardly visible on profile 1-6 (Fig. 2) whereas a higher
230 (~300 nT) amplitude anomaly is observed on profile 1-7 at the top of the Cannibal Ridge. Similarly,
231 few kilometers north of the ridge axis, a ~150 nT amplitude magnetic anomaly that is recorded on
232 profile 1-6 is absent from profile 1-7. Only one anomaly previously picked as anomaly C2A (Sauter
233 et al., 2008) and located on top of the first ridge south of the axis is continuous between the two
234 profiles (Fig. 2). On the inverted magnetization profiles (Fig. 3 and 4), an area of high magnetization
235 (up to 10 A/m) is located on the north flank of the Cannibal Ridge and is identified as the axial
236 magnetic high (profile 1-7). To the east, on profile 1-6, the same feature is shifted northward to the
237 deeper part of the axial valley. Off-axis, the magnetization is weak and associated with smooth
238 magnetic contacts.

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5.2 THE EASTERN CORRIDOR

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5.3 MAGNETIC STRUCTURE OF THE DIFFERENT TYPES OF SEAFLOOR

The eastern corridor (profiles 2-1, 2-2 and 2-3) shows a more complex morphological structure. It is at the transition from an exclusively volcanic seafloor in the west to a wide exhumed mantle domain in the east. It is characterized by a series of broad rounded serpentized peridotite ridges south of the axial valley, whereas a shallower and flatter topography prevails to the north.

The northern end of the survey (near anomaly C5, north of 27.37°S, Fig. 2, 3 and 4) shows two corrugated surfaces where the recovery of more frequent gabbroic rocks (Sauter et al., 2013) associated with a low RMBA (<20mGal, Sauter et al., 2008) suggest more robust magmatic activity.

Apart from this particular area and some thin (less than 300 m thick), small volcanic patches observed within the axial domain and at the top of some serpentized peridotite ridges, the eastern corridor is formed almost exclusively of smooth exhumed mantle surfaces associated with very little magmatic supply. Moreover, evidence was found that the ~2400 m high, 25° south dipping northern axial valley wall corresponds to the footwall of a recent large detachment fault cutting the earlier sedimented smooth inner floor and accommodating the plate separation (Sauter et al., 2013).

What has been interpreted as the central magnetic anomaly (Sauter et al., 2008) goes from a very low magnetic anomaly (<100 nT amplitude) above the detachment footwall in the west (profiles 2-3 and 2-2) to a slightly stronger anomaly ~250 nT in the deeper part of the axial valley to the east (profile 2-1; Fig. 2 and 3). Similarly, on the south flank, the anomaly picked as C2A on profile 2-1 is shifted 10 km north on profile 2-3 and is almost missing from the profile 2-2. On the conjugate plate to the north, in between the ridge axis and anomaly C5, the magnetic signal is also flat with no clear seafloor spreading anomalies and no lateral continuity; only anomaly C5 seems resolvable and quite continuous. The inverted magnetization profiles (Fig. 3) show a similar pattern to those from the western corridor: a very flat magnetization associated with smooth magnetic contacts over the exhumed mantle areas. Only anomaly C5 and very local magnetization highs, such as north of the ridge axis on profile 2-3, are observed.

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273 At the ridge axis, the magnetization of the exhumed mantle is generally low ($< 5\text{ A/m}$), but it
274 can be locally significant (e.g. up to 10 A/m on profile 1-7) and shows ill-defined magnetic contrasts
275 compared to the volcanic areas. No clear wide central block is observed in the western corridor as
276 large magnetized blocks are alternatively observed above the Cannibal Ridge (profile 1-7; Fig. 4) or
277 in the deeper part of the axial valley (profile 1-6; Fig. 4). Similarly, the central block is virtually
278 absent within the eastern corridor; a small anomaly with slightly larger magnetization (up to 8 A/m)
279 is shifting from the southern (profile 2-1; Fig. 4) to the northern axial valley wall (profile 2-3; Fig.
280 4). Off-axis, the exhumed mantle surfaces show no evidence for volcanic material (e.g. north side of
281 the profile 2-2; Fig. 4) and are characterized by low magnetizations (mostly $< 2\text{ A/m}$) without any
282 clear continuous magnetic anomaly from one profile to the other.

283 Apart from profile 2-5 showing large amplitude magnetizations, locally higher magnetization
284 cannot be associated with volcanic seafloor both at the axis and on the flanks. The presence of
285 extrusive rocks may, in some places, account for a higher magnetization but there is no unequivocal
286 link. For instance, although a lava flow is identified just north of the Cannibal Ridge on both
287 profiles 1-7 and 1-6 (Fig. 2), larger magnetization values are only observed on the eastern profile
288 (profile 1-6; Fig. 3 and 4). Similarly, although relatively higher magnetizations (up to 10 A/m) may
289 be related to the proximity of the small volcanic patch north of the axial valley wall (profile 2-3), the
290 few volcanic patches observed south of the axis of the western corridor do not produce any
291 significant magnetization ($< \pm 2\text{ A/m}$ along the profiles 2-2 and 2-3, Fig. 4)

292 The corrugated surfaces observed at the northern end of the profiles 2-2 and 2-3 are associated
293 with stronger magnetizations (up to 10 A/m) and a slightly continuous magnetic anomaly identified
294 as the anomaly C5.

295 **6. MAGNETIC PROPERTIES OF THE DREDGED SAMPLES**

296 In order to have a better understanding of the magnetic behavior of the different rock bodies in
297 the area natural remanent magnetization (NRM) and magnetic susceptibility (K) were measured on

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302 12 basalt, 29 peridotite and 10 gabbro samples dredged in the two survey areas (Fig. 1 and Table 1).
303 The magnetic susceptibility is an indication of the ability of a rock sample to acquired an induce
304 magnetization whereas NRM is a direct measurement of thermoremanence. The susceptibility is
305 plotted vs the remanence for all measured samples (Fig. 5a). Two distinct trends are observed, a
306 basalt trend with samples having low susceptibilities even when they exhibits high NRMs and a
307 peridotite trend regrouping samples having susceptibilities increasing significantly with increasing
308 NRMs (Fig. 5a) Indeed some dredged peridotites have a NRM comparable to basaltic rocks (up to 8
309 A/m), but for peridotites these high values are always associated with high susceptibility (K up to
310 0.07 SI). Gabbros tend to fall in the peridotite trend with some exception. The Koenigsberger ratio is
311 expressed as $Q=NRM/K*H$ (with H the magnetic field strength at the site) and is indicative of the
312 balance between remanent vs induced magnetization for each samples. Results are plotted on Figure
313 5b, lower Koenigsberger ratios are observed for peridotites and gabbros, half of these samples have
314 a Koenigsberger ratio below 1 indicating their magnetization is dominantly induced, this result is in
315 agreement with previous results at the MAR (Oufi et al., 2002). Basalts show ratio always above 1,
316 with a mean of ~40, indicating the strong dominance of remanent over induced magnetization.
317 Beyond these sharp, lithological based, magnetic differences it is difficult to draw any finer scale
318 magnetic behavior which could be linked to the magnetic profile. First, both NRMs and
319 susceptibilities are highly variable even for a set of samples with the same lithology collected within
320 the same dredge (see Table 1). Features in the magnetic profile are thus not easily related to rock
321 magnetic measurements. As an example the strongest magnetized peridotite samples (i.e. with the
322 higher total magnetization) were recovered within a short lateral distance of each other in the middle
323 of the north side of profiles 2-2 and 2-3 (Fig. 3). However, such a high magnetization is not
324 recorded by the deep-tow magnetic data (Fig. 3), suggesting that either the magnetized source layer
325 is thin or that high magnetizations occur only punctually, suggesting that such magnetization are
326 limited to small areas that cannot be detected by the deep-tow magnetometer.

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331 7. FORWARD MODELING

332 Seafloor spreading magnetic anomalies were identified along sea surface profiles over
333 domains that lack a volcanic upper crustal layer in the easternmost part of the SWIR (Sauter et al.,
334 2008) suggesting that other sources may play a significant role in preserving the Earth magnetic
335 field polarity. Along both the MAR and the SWIR serpentinized peridotites have been suspected to
336 carry “more positive” magnetization amplitudes in areas of thin crust (Hosford et al., 2003;
337 Tucholke et al., 2008). This observation is supported by the Koenigsberger ratio of the SWIR
338 serpentinized peridotites which is often less than 1. This shows that, unlike basaltic rocks, induced
339 magnetization may be significant for these rocks. We therefore perform two different forward
340 modeling, one based on a 500 m thick basaltic layer with a dominantly remanent magnetization thus
341 preserving the Earth magnetic field polarity and another based on an induced magnetized layer, in
342 order to test the contribution of volcanic rocks versus serpentinized peridotites. We disregard the
343 contribution of a lower crustal layer made of gabbroic rocks that is volumetrically scarce in the
344 samples dredged within the exhumed mantle domains.

345 7.1 SEAFLOOR SPREADING MODEL

346 The seafloor spreading model calibrated on the volcanic seafloor ([profile 2-5](#)) was compared
347 to the magnetic profiles acquired above the exhumed mantle domains of both eastern and western
348 corridors. In the absence of a high amplitude central anomaly, the axial Brunhes block was centered
349 either at the bathymetric axis or underneath the central magnetic anomaly. Fig. 2 shows the
350 predicted magnetic anomaly along each across-axis profile (dashed black line). The parameters
351 derived from profile 2-5 give a poor fit to the observed magnetic field for the two corridors. The
352 picked axial anomaly and anomaly C2A or C3A on sea surface magnetic profiles (Cannat et al.,
353 2006) are not clearly observed on the deep-tow profiles over the exhumed mantle domains.
354 Furthermore, the modeled anomaly C5 at the end of both profiles 2-2 and 2-3 appears to be shifted a
355 few kilometers to the north with respect to the previously picked anomaly C5. This offset may be
356 explained by either asymmetrical spreading, changes in spreading rate between [anomalies](#) C3A and

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359 C5 or a mislocation of the central Brunhes anomaly.

360 7.2 INDUCED MAGNETIZED MODEL

361 To account for the unconstrained lateral and vertical variations in both intensity and direction
362 of the remanent component of magnetization, we tested whether a uniform induced magnetized
363 layer could solely account for the observed magnetic anomalies. As unaltered peridotites have a very
364 weak susceptibility the depth extent of serpentinization has to be determined. Seismic velocities in
365 the exhumed mantle domains of the Iberian margin, as well as over slow-spreading ridges, suggest a
366 high serpentinization degree (greater than 75%) in the first 2 km below the seafloor (Minshull et al.,
367 1998; Chian et al., 1999; Dean et al., 2008). Thus, based on average NRM values measured on our
368 dredge samples but also on serpentinites at ODP Holes 897D, 899B, 1070A and 1277 at the Iberia-
369 Newfoundland margins (Zhao et al., 2001), we use a 1.5 A/m constant magnetization and 2 km
370 constant thickness draped on the bathymetry as a source for the magnetic anomalies. In such models
371 we assume that the whole magnetic signal is exclusively related to the seafloor topography.

372 The results are represented in Figure 2 by the thin black continuous lines. In the western
373 corridor, there is a poor fit between the synthetic and the observed magnetic anomaly along profile
374 1-5 and especially in the axial valley of profile 1-6. However, the agreement is better on profile 1-7,
375 even within the central domain. The only continuous anomaly between both profiles 1-6 and 1-7,
376 picked as anomaly C2A by Sauter et al. (2008), is well marked on both profiles. In the eastern
377 corridors, where the areas of flat topography correspond to a flat observed magnetic field, the data
378 are slightly comparable to the model except in the axial domain. Over the volcanic seafloor, the
379 synthetic magnetic anomaly fits poorly the observed magnetic anomaly along the profile 2-5 in both
380 axis and off-axis regions.

381 8. DEPTH OF THE MAGNETIC SOURCES

382 Usually, sources for marine magnetic anomalies are considered to be located in the extrusive

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384 upper part of the oceanic crust within a layer of constant thickness. Following Bronner et al. (2013),
385 we use the equivalent layer method for the estimation of the thickness of this layer. Dampney (1969)
386 has shown that the equivalent layer should be located within a certain range of depth below the
387 measurement surface for avoiding both the aliasing effect in the computed field and an ill-
388 conditioned inversion matrix. In our case, we use the top seafloor as an upper bound because the
389 altitude of the TOBI largely exceeds the data spacing (~10 m) and we iteratively increase the depth
390 of the equivalent layer until (1) the loss of short wavelength in both computed field and
391 magnetization, and/or (2) the appearance of oscillations in the solutions. Indeed, an equivalent
392 source layer located too far below the measurement surface makes the matrix representing the
393 distance between the equivalent sources and the observation points ill conditioned and the associated
394 solution unreliable (Dampney, 1969). Therefore, we assume that this lower bound provides a first
395 indication relative to the maximum depth of the “true” causative sources (i.e. the source layer
396 thickness).

397 Applying this method to our survey, in the case of the “volcanic” profile ([profile 2-5](#)), the
398 whole frequency content of the measured field was well retrieved for dipoles located around 500 m
399 below the seafloor (Fig. 6). Shallower and deeper solutions lead respectively to the appearance of
400 high frequency oscillations in the computed field and loss of resolution in the computed field and
401 magnetization solution. At the opposite, for the profiles acquired above the exhumed mantle
402 domains (e.g. profile 2-2; Fig. 6), the weakness of the signal associated with a quasi absence of short
403 wavelength anomalies allow reasonable solutions for both synthetic field and magnetization within a
404 wider range of depth (up to 2000 m below the seafloor; Fig. 6).

405 As for any methods used to estimate the depth of magnetic sources, these results have to be
406 taken with care. The maximum depth of 500 m obtained above the volcanic seafloor is in agreement
407 with the 500 m thick basaltic layer generally used to account for the marine magnetic anomalies at
408 mid oceanic ridges. This suggests that the 2D hypothesis used here is reliable in the case of a 2D
409 homogeneous crustal accretion but we do not have much constraints on the magnetization structure

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111 of the exhumed mantle domains and the 2D assumption may lead to some errors. The deeper
112 solutions found for the sources in exhumed mantle domains mainly suggest that the short
113 wavelength magnetic anomalies recorded above the volcanic areas are missing from the profiles
114 acquired above serpentinized peridotites. This can be explained by deeper sources or by smoother
115 magnetic contacts.

116 9. DISCUSSION

117 9.1 SEAFLOOR SPREADING ANOMALIES

118 Because the shape of the marine magnetic anomalies strongly depends on the distance
119 between two polarity reversals (i.e. frequency of polarity reversal versus spreading rate), the
120 magnetic reversal pattern along ultra-slow spreading ridges is often blurred. At the SWIR,
121 geomagnetic reversals used for the resolution of spreading rates for the last 24 Ma have often been
122 restricted to long reversals of constant polarity, such as chrons C5 and C6, of about 1 Myr duration
123 (Patriat et al., 2008). Moreover, in the eastern part of the SWIR, large variations in both crustal
124 thickness (inferred from RMBA; Cannat et al., 2006) and lithology (from volcanic basalt to
125 tectonized ~~serpentinized~~ peridotites) are associated with different modes of seafloor generation
126 (volcanic extrusion ~~or~~ mantle exhumation). In addition to the ultra-slow spreading rate, these
127 accretion modes are responsible for the complexity of the magnetic signal in this area. Short
128 reversals such as C3 or C2A (~0,5 Myr) are detectable only above thick volcanic crust associated
129 with minor tectonic activity. Therefore, the identification of magnetic anomalies from sea surface
130 magnetic profiles above smooth seafloor was mainly extrapolated from the surrounding volcanic
131 areas (Sauter et al., 2008). However, the comparison of mapping of exhumed mantle domains from
132 multibeam bathymetric data at 150 m resolution and from TOBI images at <10 m resolution reveals
133 that sources of some of these magnetic anomalies identified along sea surface profiles were
134 erroneously attributed to volcanic seafloor by Sauter et al. (2008) and may thus not be related to
135 polarity changes of the Earth magnetic field.

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140 We are now able, with the high-resolution deep-tow data, to provide a precise analysis of the
141 magnetic signal with the respect to the geological nature of the seafloor. We confirm, as observed by
142 Searle and Bralee (2007) in this region, that despite an ultra-slow spreading rate, marine magnetic
143 anomalies are still well identifiable above volcanic seafloor. However, these identifications are
144 much more difficult above the exhumed mantle domains where the magnetic pattern is highly
145 variable from one profile to another. This is well illustrated in our studied corridors where closely
146 spaced magnetic profiles show a very heterogeneous magnetic signal. Except for anomaly C5, a
147 simple seafloor-spreading model (calibrated on ~~profile 2-5~~) does not fit the magnetic anomaly
148 pattern, even for the central anomaly (Fig. 2). Moreover, although the exhumed mantle domains are
149 expected to be formed by asymmetrical detachment faulting, there is no evidence for lateral
150 discontinuity between exhumed mantle areas and symmetrically accreted volcanic crust for which
151 ages of accretion are quite well constrained (e.g. Chrons younger than C5; Searle and Bralee, 2007).
152 We thus suggest that both mantle exhumation and volcanic accretion are almost contemporary
153 leading to a reasonable lateral continuity (in terms of age) between different types of seafloor.

154 9.2 CONTRIBUTION OF MANTLE DERIVED ROCKS

155 The ferromagnetic behavior of serpentinized peridotites has been shown to be directly linked
156 to the serpentinization process (Dunlop and Prévot, 1982). Magnetite is formed during
157 serpentinization from the interaction between fluids and ferromagnesian minerals such as olivine
158 and pyroxene. It has been suggested that a high degree of serpentinization (above ~75%) is
159 necessary for the acquisition of both significant susceptibility and NRM (Oufi et al., 2002).
160 However, highly variable NRMs are observed from one ODP site to another and also between
161 samples drilled in a single ODP Hole (Oufi et al., 2002). In our study, a similar magnetic behavior is
162 observed for the dredged peridotites. Values of NRM and susceptibility can be significant but are
163 highly heterogeneous. The susceptibilities and NRM values in our dredged samples fall in the lower
164 range of values reported for drilled abyssal peridotites (Oufi et al., 2002). Nevertheless, our samples
165 show similar Koenigsberger ratios to these of drilled peridotites and ~~they~~ are strictly inferior to those
166 expected for the extrusive upper layer of the oceanic crust (e.g. ~50 to 300; Marshall and Cox,

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168 1971). Although peridotites outcropping at the seafloor are subject to low-temperature alteration, we
169 consider that the magnetic properties of the dredged peridotites are representative of a magnetic
170 source layer in the exhumed mantle domains. Based on our data, we suggest that the high variability
171 in NRM intensity combined with the large range of susceptibility and low Koenigsberger ratio make
172 this layer of serpentinized peridotite magnetically weak and variable. This is in agreement with a
173 recent study of serpentinized peridotites samples from both the Mirdita ophiolite and ocean drilling
174 sites showing that strongly serpentinized (> ~60%) peridotites have variable Königsberger ratios and
175 are affected by randomly oriented, low stability components that result in incoherent NRM
176 directions at both the site and regional scales (Maffione et al., 2014).

177 The presence of such layer is confirmed by the high-resolution deep-tow profiles, which
178 display highly variable magnetic pattern from one profile to another, even where they are closely
179 spaced. The magnetic anomalies observed above exhumed mantle are weak (<100 nT), and lack
180 short wavelength anomalies, suggesting deeper magnetic sources than in the volcanic seafloor.
181 Although it is possible to reproduce some magnetic pattern using a induced magnetic layer, the
182 whole magnetic signal is not retrieved, especially within the axial domain. Similarly, the spreading
183 model does not fit the observed data. It is thus likely that the sources combine both induced and
184 remanent magnetization and vary from one local area to another.

185 Furthermore, in the western and more magmatic part of the SWIR (54-56°E), a significant
186 along-axis decrease in magnetization produces the disappearance of the magnetic reversal patterns
187 in the deepest parts of ridge discontinuities (Sauter et al., 2004). This observation was linked to the
188 thinning of the upper part of the oceanic crust due to a decreasing magmatic budget toward the
189 segment ends. It has also been shown, at the MAR (13-15°N), within a highly complex accretion
190 context combining detachment faulting and freshly erupted seafloor, that the magnetic pattern could
191 be significant in amplitude but highly heterogeneous on a scale of ~5 km (Mallows and Searle,
192 2012) leading to difficulties in the identification of the spreading anomalies, even for the large
193 Brunhes central block. This further suggests that a sufficiently homogeneous upper crust is required
194 to produce well marked marine magnetic anomalies, and that in the absence of this main magnetic

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499 source made of extrusive (and perhaps intrusive material), exhumed serpentinized peridotites are not
500 sufficiently uniform magnetic sources to produce undisputable seafloor spreading magnetic
501 anomalies.

502 9.3 CORRUGATED SEAFLOOR AND THE MAGNETIC SIGNAL

503 A corrugated surface is observed below the identified anomaly C5 toward the north end of
504 profiles 2-2 and 2-3 in the eastern corridor. This area is strongly magnetized (up to 10 A/m; Fig. 2
505 and 3) and displays high amplitude magnetic anomaly up to 450 nT. This corrugated surface is
506 surrounded by lineated volcanic terrains and the conjugate anomaly C5, on the opposite flank, was
507 identified over well established volcanic crust. This observation, together with a frequent recovery
508 of gabbro in this area (Sauter et al., 2013), suggests that the anomaly C5 was emplaced in a more
509 robust magmatic accretion context, before or just at the onset of continuous mantle exhumation. We
510 thus speculate that in this particular area the magnetization is carried by extrusive or intrusive
511 material rather than by peridotites.

512 9.4 VOLCANIC SEAFLOOR AND THE MAGNETIC SIGNAL

513 Some small volcanic patches have been identified respectively just north of the axial valley of
514 the western corridor and within the axial region of the eastern corridor (Fig. 4). It is not clear
515 whether this extrusive material always accounts for higher magnetization than over the smooth,
516 exhumed mantle seafloor. The magnetic data rather confirm the interpretation of Sauter et al. (2013)
517 based on deep-tow sonar images that these volcanics are very thin and discontinuous flows, not
518 exceeding a hundred meters of thickness, and thus do not correspond to large enough sources to be
519 identified by the deep-tow magnetometer.

520 9.5 MARINE MAGNETIC ANOMALIES AT OCEAN-CONTINENT TRANSITION

521 Mantle exhumation is one of the proposed mechanisms responsible for the formation of the
522 transitional domains at magma-poor rifted continental margins where serpentinized mantle-derived
523 rocks have been drilled (Tucholke and Sibuet, 2007). Sibuet et al. (2007) proposed that a strong

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526 magnetization (up to 9 A/m) can be produced by the serpentinization of a 2 to 3 km thick fractured
527 layer, within the root of an active detachment fault at an embryonic spreading center. Based on
528 NRM intensity measured in ODP holes at the Iberia margin, these authors argued that this first
529 serpentinization phase is sufficient to preserve the polarity of the ambient magnetic field. They
530 suggest that only the upper ten meters below the seafloor are affected by cold-water alteration that
531 produces incoherent magnetic properties. Like on the SWIR, the exhumed mantle domains of the
532 Iberia-Newfoundland margins are characterized by a weak and ill-defined magnetic signal. At these
533 margins, only the seaward termination of the exhumed mantle domain is associated with a slightly
534 linear and high amplitude (up to 1000 nT) magnetic anomaly (“J anomaly”). This anomaly was
535 interpreted as the end of the M sequence of spreading anomalies and its amplitude was explained by
536 a strongly serpentinized crust (Srivastava et al., 2000).

537 No clear seafloor spreading anomaly is observed over the exhumed mantle areas of the SWIR,
538 neither where active detachment faulting is identified nor on the flanks. This leads to the conclusion
539 that the serpentinization process is not sufficiently homogeneous to produce stable large remanent
540 magnetization. We suggest that the heterogeneous magnetization of the serpentinized peridotites is
541 strongly depending on the fluid-rock interactions, the temperature, the mineral composition and the
542 tectonic context. Therefore, in view of the low magnetization of the young (<11Ma) serpentinized
543 rock at the SWIR, it is unlikely that strong magnetic anomalies could be related solely to
544 serpentinization; this would be even more true at >100 Ma old OCTs. This instead supports the
545 hypotheses that (1) intrusive or extrusive material is required (Bronner et al., 2011; Russell and
546 Whitmarsh, 2003) to account for a significant magnetic signal in the exhumed mantle domains of
547 OCTs and that (2) the interpretation of this signal as resulting from seafloor spreading is precluded
548 in the absence of a homogeneous and well established upper oceanic crust. Consequently, the
549 kinematic reconstructions of magma-poor passive margins using weak anomalies identified over
550 exhumed mantle domains need to be taken with caution.

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553 **10. CONCLUSION**

554 We have investigated the magnetic structure of newly discovered large exhumed mantle
555 domains of the SWIR (Sauter et al., 2013) combining high resolution side-scan sonar images, deep-
556 tow magnetic data and results from dredge sampling. We show that the seafloor spreading magnetic
557 pattern disappears from the volcanic seafloor toward the exhumed mantle domains. Forward
558 modeling allows a reasonable fit to the observed magnetic anomalies over the volcanic seafloor.
559 However, the lack of a central magnetic anomaly and the highly heterogeneous and weak magnetic
560 pattern observed above exhumed mantle-derived rocks prevents any identification of polarity
561 reversals. Moreover, analysis of the magnetic properties of the dredge samples shows that
562 serpentinized peridotites as well as gabbros are highly variable magnetic sources. We conclude that
563 the serpentinization process is not sufficiently homogeneous to produce a significant stable
564 magnetization at the scale of the exhumed mantle domains of the SWIR and that serpentinized
565 peridotites are not able to contribute to regionally-coherent patterns of oceanic magnetic anomalies.
566 We further suggest that a homogene~~ous~~ volcanic upper crust associated with minor tectonic activity
567 is required to record well-defined seafloor spreading magnetic lineations.

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569 **Figure 1** : Bathymetric map with the location of the two survey areas (western and eastern corridor).
570 Magnetic anomaly picks are from Sauter et al. (2008). The nature of the seafloor was deduced either
571 from the side-scan images when available (Sauter et al. 2013) or from the multibeam bathymetric
572 data (Cannat et al. 2006). The dredges numbers and the proportion of rocks by weight shown as pie
573 charts are from Sauter et al. (2013). We have only shown the dredges for which we have measured
574 the magnetic properties (see table 1 and Fig. 3 and 5)

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575 **Figure 2** : 2D magnetic profiles recorded within the two survey areas along the profiles shown in
576 Fig. 1. Magnetic data (continuous red lines) have been upward continued to an altitude of 1200 m
577 below the sea level. Broken lines correspond to the magnetic anomaly predicted by a 14 mm/a.
578 seafloor spreading model calibrated on the volcanic seafloor (Profile 2-5) with a 500m thick source
579

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584 layer and a 10 and 5 A/m magnetization for the axial and off axis blocks, respectively. Black solid
585 lines correspond to a model based on a 2 km thick source layer for which a solely induced and
586 uniform magnetization is applied (1.5 A/m). Interpretations from the side-scan images are shown
587 below the bathymetry profiles (from Sauter et al., 2013). The vertical grey area indicates the
588 location of the axial valley.

589 |
590 **Figure 3 :** Inverted magnetization for the western (a) and eastern survey area (b). Colored strips
591 show the calculated magnetization values along the magnetic anomaly profiles (black lines). Shaded
592 relief images are shown in background. Red circles are sized relatively to the NRM values of
593 dredged basalts whereas green circles correspond to NRM values measured on dredged peridotites
594 (see table 1). The dredge number is shown in the white box near the circles. The thin black lines
595 corresponding to the edges of the volcanic seafloor are from Sauter et al. (2013). Picking of
596 magnetic anomalies is the same as Figure 1. 14 mm/a reversal pattern is superimposed for
597 comparison in fig.3b.

599 **Figure 4 :** 3D bathymetric view of the two survey areas. The inverted magnetization (colored strips)
600 and the edges (from Sauter et al., 2013) of both the corrugated surfaces (purple lines) and the
601 volcanic seafloor (white faded areas) are draped on the multibeam bathymetric map. The
602 magnetization scale is the same for the two survey areas. The black lines indicate the edges of the
603 TOBI side-scan swath.

605 **Figure 5 :** (a) Natural remanent magnetization (NRM) from dredged peridotites, basalts and gabbros
606 as a function of the magnetic susceptibility (K). (b) Koenigsberger ratio (Q) for serpentinized
607 peridotites (SP), basalts (B) and gabbros (G). Note that Q has a logarithmic scale.

609 **Figure 6 :** Comparison between the deep-tow observed magnetic field, along the profiles 2-2 and 2-
610 5, and the magnetic computed field along the TOBI path for different depths of inferred magnetized
611 dipoles from 0 to 2000 m below the seafloor. In the lower panel, the magnetization solution is also

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525 [reported along the same profiles for different depths](#). The shallowest and deepest dipoles lead
526 respectively to the appearance of high frequency oscillations in the computed field and loss of
527 resolution in both the computed field and the magnetization solution. The best compromise is found
528 for dipoles located around 500 m below the seafloor for the profile [2-5](#) acquired above the volcanic
529 crust and 1000 m for the profile [2-2](#) collected above the exhumed mantle derived rocks. A
530 significant loss in resolution is observed for dipoles located below 1000 m in the case of the
531 volcanic crust whereas both the computed field and the magnetization solution are quite well
532 preserved up to 2000 m for the case of exhumed mantle seafloor.

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533 [Acknowledgements](#)

534 [We thank officers and crew of the R/V Marion Dufresne for their assistance during the](#)
535 [“Smoothseafloor” cruise. We also thank anonymous reviewers for their constructive comments that](#)
536 [significantly improved this manuscript. Funding was provided by ANR grant ‘Rift2Ridge’ and](#)
537 [support by INSU/CNRS and IPEV.](#)

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