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Synchrotron FTIR imaging of OH in quartz mylonites

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66 **Abstract**

67

68 Methods of measuring OH absorption bands of fluid inclusions and hydrogen
69 defects in deformed quartz-rich rocks at high spatial resolution are described,
70 using synchrotron infrared IR radiation coupled with a Fourier transform infrared
71 FTIR microscope, and applied to imaging OH in mylonites from the footwall of
72 the Moine Thrust (from the Stack of Glencoul, NW Scotland Caledonides) and
73 hanging wall of the Main Central Thrust (Sutlej Valley, Himalaya of NW India).
74 Previous measurements of water in deformed quartzites using conventional FTIR
75 instruments, through apertures of 50-100 μm for specimens $\sim 100 \mu\text{m}$ in thickness
76 have shown that water contents of larger grains vary from one grain to another.
77 However, the non-equilibrium variations in water content between neighboring
78 grains and within quartz grains cannot be interrogated further without greater
79 measurement resolution, nor can water contents be measured in finely
80 recrystallized grains without including absorption bands due to fluid inclusions,
81 films, and secondary minerals at grain boundaries.

82 Synchrotron IR radiation is brighter and more collimated than offered by
83 conventional FTIR global light sources, and we have been able to distinguish and
84 measure OH bands due to fluid inclusions, hydrogen point defects, and secondary
85 hydrous mineral inclusions through an aperture of 10 μm for specimens $> 40 \mu\text{m}$
86 thick. Doubly polished IR plates can be prepared with thicknesses down to 4-8
87 μm , but measurement of small OH bands is currently limited by strong
88 interference fringes for samples $< 25 \mu\text{m}$ thick, precluding measurements of water
89 within individual, finely recrystallized grains. By translating specimens under the
90 10 μm IR beam by steps of 10 to 50 μm , using a software-controlled x-y stage,
91 spectra have been collected over specimen areas of nearly 4.5 mm^2 . Integrated
92 absorbance of the broad OH band at 3400 cm^{-1} due to fluid inclusions was mapped
93 and water content contoured, after subtracting the 3600 cm^{-1} band due to micas.



94 Maps of the 3600 cm^{-1} band absorbance show concentrations of micas at grain
95 grain boundaries, as well as finely dispersed mica inclusions within quartz grains.
96 Mylonitic Cambrian quartzites deformed under greenschist facies
97 conditions in the footwall to the Moine Thrust exhibit a large and variable 3400
98 cm^{-1} band absorbance, and maps of water content corresponding to fluid inclusions
99 show that inclusion densities correlate with deformation and recrystallization
100 microstructures. Water contents of quartz grains in these mylonites are
101 comparable to those of wet quartz varieties deformed in laboratory experiments.
102 OH absorption bands of large clear quartz grains of mylonitic orthogneisses and
103 paragneisses deformed under lower-upper amphibolite facies conditions in the
104 hanging wall to the Main Central Thrust are much smaller, and dominated by
105 sharp bands at 3595 cm^{-1} to 3379 cm^{-1} due to hydrogen point defects that appear to
106 have uniform, equilibrium concentrations in the driest samples. The broad OH
107 band at 3400 cm^{-1} in these rocks is much less common. The lack of detectable
108 water in highly sheared Main Central Thrust mylonites offers a challenge to our
109 understanding of quartz rheology. However, where water absorption bands can be
110 detected and compared with deformation microstructures, OH concentration maps
111 provide information on the histories of deformation and recovery, and evidence for
112 the introduction and loss of fluid inclusions, and water weakening processes.

113

114

115 *Keywords:* FTIR Mapping, Water weakening, Dislocation Creep, Quartz

116

116 **1 Introduction**

117

118 Quartz mylonites sheared at middle to lower levels of the continental crust exhibit
119 microstructural and textural evidence of dislocation creep, a process that is widely
120 believed to require water weakening in framework silicates. The effects of water
121 on dislocation creep of quartz, including the nucleation, glide, climb and recovery
122 of dislocations, and recrystallization are well known from: 1) experimental studies
123 of natural crystals, in which water was introduced into grain interiors (e.g., Griggs,
124 1967; Blacic, 1975, 1981; FitzGerald et al., 1991), 2) studies of synthetic and
125 natural quartz varieties with large initial water contents (e.g., Griggs and Blacic,
126 1965; Hobbs, 1968; Baeta and Ashby, 1970; Kekulawala et al., 1978; Kirby and
127 McCormick, 1979; McLaren et al., 1983; Linker et al., 1984; Gerretsen et al.,
128 1989; Muto et al., 2011; Holyoke and Kronenberg, 2013; Stünitz et al., in press),
129 and 3) quartzites and polycrystalline quartz aggregates with water added or
130 removed before or during experiments (e.g., Jaoul et al., 1984; Kronenberg and
131 Tullis, 1984; Tullis and Yund, 1989; Hirth and Tullis, 1992; Gleason and Tullis,
132 1995; Post et al., 1996; Chernak et al., 2009). In addition to the parametric
133 approach taken to study water weakening in deformation experiments, key insights
134 into water weakening come from infrared IR spectroscopy of quartz samples that
135 have been deformed in the lab under varying hydrothermal conditions, explicitly
136 adding or removing water or changing the fugacity of H₂O at fixed temperature
137 and/or strain rate. OH absorption bands characteristic of different hydrogen
138 defects and forms of molecular water can be identified and their size used to
139 determine concentration (e.g., Kats, 1962; Griggs and Blacic, 1965; Aines and
140 Rossman, 1984; Aines et al., 1984; Stipp et al., 2006).

141 Water weakening in the continental crust is inferred because of the high
142 laboratory strengths exhibited by quartz and feldspars in the absence of water
143 (Griggs, 1967; Heard and Carter, 1968; Tullis and Yund, 1977, 1980; Tullis, 1983;
144 Blacic and Christie, 1984) and the postulated effects of water on point defects and



145 disruptions of fully linked Si-O bonds (Griggs, 1974; Hirsch, 1979; Hobbs, 1981;
146 Paterson, 1989). With the advent of Fourier transform infrared FTIR spectroscopy
147 and IR microscopes, water and hydrogen defects in naturally deformed quartz
148 have been reported with OH contents of 300 ppm to >10,000 ppm (molar H/10⁶ Si,
149 Kronenberg and Wolf, 1990; Kronenberg et al., 1990; Nakashima et al., 1995;
150 Gleason and DeSisto, 2008; Seaman et al., 2013; Finch et al., 2016), comparable
151 to those required for water weakening in laboratory experiments. In small granitic
152 shear zones deformed at greenschist conditions, water contents of quartz grains
153 appear to correlate with finite strain (Kronenberg et al., 1990; Gleason and
154 DeSisto, 2008), and in the much larger Median Tectonic Line of Japan,
155 intragranular water contents increase towards its center (Nakashima et al., 1995).
156 However, OH concentrations of quartz of granitic rocks deformed at higher
157 temperatures can be much lower (20 - 100 ppm; Han et al., 2013; Kilian et al.,
158 2016), with IR spectra dominated by small sharp OH bands of hydrogen point
159 defects (Kilian et al., 2016) that are not thought to weaken quartz. Moreover,
160 quartz water contents have been reported that show a trend of decreasing OH
161 content towards the center of a high grade shear zone (Finch et al., 2016). Maps of
162 OH content, constructed from FTIR spectra of deformed granitic rocks (Seaman et
163 al., 2013) show compelling relationships between water content and
164 microstructures generated during deformation, suggesting that water contents are
165 reduced during recrystallization and partial melting.

166 Much as deformation microstructures and textures provide a link between
167 our understanding of deformation mechanisms activated in deformation
168 experiments at high laboratory strain rates and the mechanisms governing
169 plasticity and creep of shear zones at low natural strain rates (Snoke et al., 1998;
170 Heilbronner and Barrett, 2014), IR spectroscopy can provide a link between our
171 understanding of water weakening in the lab and in nature. In this paper, we
172 report on methods of FTIR spectroscopy to characterize OH absorption bands and
173 image OH contents in quartz (and other nominally anhydrous minerals) at higher



174 resolution than is possible using conventional instruments, coupling synchrotron
175 IR radiation with FTIR spectroscopy. We then apply these methods to mylonitic
176 quartzites in the footwall to the Moine Thrust in the NW Scotland and to mylonitic
177 quartz-rich orthogneisses and paragneisses in the hanging wall to the Main Central
178 Thrust in NW India. OH absorption bands are readily measured through an
179 aperture of 10 μm for samples $> 40 \mu\text{m}$ thick, sufficient to measure water contents
180 in highly strained quartz grains (ribbons) and less deformed porphyroclasts and
181 augen of the Moine Thrust mylonites. However, water contents of small
182 individual recrystallized grains of the Moine Thrust mylonites could not be
183 measured since spectra of doubly polished IR specimens $< 25 \text{ mm}$ thick display
184 strong interference fringes. OH absorption bands can be measured for the larger
185 grains of the Main Central Thrust mylonites, including original and recrystallized
186 grains.

187 Images of intragranular water content in Moine Thrust and Main Central
188 Thrust mylonites, based on contours of integrated absorbance of the broad OH
189 band at 3400 cm^{-1} correspond to populations of fluid inclusions, and to
190 deformation and recrystallization microstructures. Planar arrays of secondary
191 fluid inclusions within quartz grain interiors provide evidence of increases in water
192 during brittle microcracking, followed by crack healing and processes of fluid
193 inclusion evolution and redistribution. Contoured integrated absorbances of the
194 sharper band at 3600 cm^{-1} due to OH in micas show their spatial distributions at
195 quartz grain boundaries and interiors.

196

197 **2 Selected quartz mylonites**

198

199 The quartz mylonites selected to test synchrotron FTIR spectroscopy and imaging
200 of OH come from the footwall of the Moine Thrust at the Stack of Glencoul, NW
201 Scotland (Christie, 1963; Law et al., 1986, 2010) and the hanging wall of the Main
202 Central Thrust exposed in the Sutlej Valley, NW India (Law et al., 2013; Stahr,



203 2013; Law, 2014). Both of these thrust faults are orogen-scale shear zones with
204 penetrative deformation on the Moine Thrust accommodating shortening at the
205 foreland edge of the Caledonian orogeny (e.g., Peach et al., 1907; Elliot and
206 Johnson, 1980; Law et al., 1986, 2010; Butler, 2010; Law, 2010; Law and
207 Johnson, 2010; Dewey et al., 2015) and penetrative shear strains on the Main
208 Central Thrust accommodating southward directed extrusion/exhumation of the
209 overlying Greater Himalayan slab (e.g., Grujic et al., 1996; Grasemann et al.,
210 1999; Godin et al., 2006; Law et al., 2013). Mylonitic grain shape foliations are
211 well developed in rocks of both fault zones and mineral stretching lineations are
212 parallel to the fault transport directions. Deformation microstructures and textures
213 of quartz in these mylonites indicate that dislocation creep was the predominant
214 deformation mechanism involving both basal and prism slip systems, internal
215 recovery and dynamic recrystallization.

216 Structural distances of samples below the Moine Thrust and above the
217 Main Central Thrust, together with available information on deformation
218 temperatures and temperatures/pressures of metamorphism, are summarized in
219 Table 1. All samples from the footwall to the Moine Thrust are mylonitic
220 Cambrian quartzites. Samples from the hanging wall to the Main Central Thrust
221 on the NW Sulej transect are penetratively deformed orthogneisses in which
222 intensities of grain shape fabrics, traced downwards towards the thrust surface,
223 increase (associated with the transition from grain boundary migration to subgrain
224 rotation recrystallization) and dynamically recrystallized grain sizes decrease (Law
225 et al. 2013; their fig. 6). Samples from the more hinterland-positioned Eastern
226 Sulej transect are less obviously foliated paragneisses with more granular textures
227 due to extreme quartz grain boundary mobility. At a given distance above the
228 Main Central Thrust, samples from the more hinterland positioned Eastern Sulej
229 transect have much larger recrystallized grain sizes than on the NW Sulej transect
230 (compare figs. 15 and 6 in Law et al. 2013).



231 Mylonitic Cambrian quartzites in the footwall to the Moine Thrust at the
232 Stack of Glencoul (Assynt region) display highly flattened relict grains aligned
233 parallel to foliation, with grain shape aspect ratios up to 50:1 to 100:1 and smooth
234 undulatory extinction between crossed polarizers that have been described as
235 quartz ribbons (Bonney, 1883; Christie, 1960, 1963; Weathers et al., 1979; Law et
236 al., 1986). At the margins of the larger quartz grains, more equant, finely
237 recrystallized grains overprint these elongate high-strain grains, with the
238 proportion of new to old grains increasing structurally upwards towards the Moine
239 Thrust plane (Christie, 1960; Weathers et al., 1979; Law et al., 1986). Rare grains
240 of feldspar and quartz aligned in “*mechanically strong*” orientations are relatively
241 equant, and appear as augen or globular quartz grains. Quartz c-axes exhibit
242 strong lattice preferred orientations in both deformed old grains and recrystallized
243 grains, with symmetrical Type 1 (Lister, 1977) cross-girdle fabrics at distances >
244 150 mm beneath the thrust plane and increasingly asymmetric cross-girdle to
245 single girdle fabrics closer to the thrust plane (Law et al., 1986, 2010). These
246 fabrics reflect general flattening strains accommodated by quartz basal and prism
247 slip, with variations in estimated flow vorticities and partitioning of strain between
248 original and recrystallized grains (Law et al., 2010; Law 2010). Micas are highly
249 aligned parallel to foliation, with coarse muscovite grains at quartz grain
250 boundaries and fine, dispersed micas within quartz grain interiors. Optical
251 microstructures show evidence of quartz recrystallization by bulge nucleation and
252 subgrain rotation with relatively uniform mean recrystallized grain sizes of ~15
253 μm (Christie et al., 1954; Weathers et al., 1979). Microstructures imaged by
254 transmission electron microscopy TEM show dense arrays of curved free
255 dislocations, subgrain walls, and fine fluid inclusions that decorate dislocations
256 (Weathers et al., 1979; Ord and Christie, 1984). Earlier FTIR measurements of
257 Moine Thrust mylonites from a number of locations in the Assynt region revealed
258 large OH absorption bands characteristic of milky quartz with OH contents of
259 1500 to 7500 ppm (Kronenberg and Wolf, 1990).



260 Quartz-rich Greater Himalayan Series orthogneisses and paragneisses in
261 the hanging wall to the Main Central Thrust exposed in NW and Eastern Sutlej
262 transects, respectively, include quartz mylonites, quartz-mica schists and quartz-
263 garnet schists (Vannay and Grasemann, 1998; Grasemann et al., 1999; Law et al.,
264 2013). Quartz grain shapes are not as highly elongate as observed in the Moine
265 Thrust mylonites, owing to their extensive recrystallization, with mean grain sizes
266 on the NW transect that vary with structural level from 200-250 μm (grain
267 boundary migration microstructures) at ~ 1000 m above the Main Central Thrust to
268 75-95 μm at 200-750 m above the thrust, and 35-60 μm (dominantly subgrain
269 rotation microstructures) at ~ 75 m above the thrust surface (Law et al., 2013). As
270 noted above, quartz recrystallized grain sizes are much larger at a given structural
271 height above the thrust plane on the hinterland-positioned Eastern transect with
272 grain sizes commonly ranging from 250 - >1000 μm (0.25-1.0 mm) at ~ 25 m
273 above the thrust to >1 -2 mm at 200 m and higher above the thrust. Quartz grain
274 interiors on the Eastern transect display mild undulatory extinction with highly
275 aligned subgrain boundaries that give the appearance of chessboard extinction and
276 irregular, non-planar grain boundaries that envelope neighboring mica grains,
277 suggestive of high grain boundary mobilities at the time of peak metamorphism
278 and deformation. Quartz *c*-axes show strong lattice preferred orientations with
279 symmetrical and slightly asymmetrical cross-girdle fabrics on both transects,
280 providing evidence of simultaneous basal and prism slip during plane strain
281 deformation, with varying amounts of pure shear and simple shear (Law et al.,
282 2013). Coarse-grained muscovite and biotite are highly aligned parallel to
283 foliation, both at the boundaries of quartz grains and within quartz grain interiors.

284 While both the Moine Thrust and the Main Central Thrust mylonites were
285 deformed by dislocation creep, their deformation temperatures were very different
286 and the mechanisms of recrystallization, accommodating dislocation creep differ
287 accordingly. Deformation temperatures for the Moine Thrust footwall mylonites
288 at the Stack of Glencoul are estimated at between 300° and 350°C based on illite



289 crystallinity (Johnson et al., 1985). Opening angles of *c*-axis fabrics measured
290 separately on old and recrystallized grains (Law et al., 2010) indicate far higher
291 apparent deformation temperatures using the Kruhl (1998) fabric opening angles
292 thermometer (390°-440°C and 475°-530°C, respectively) than are compatible with
293 their prehnite-pumpellyite to lower greenschist (chlorite) facies phyllosilicates
294 (Law, 2014). Deformation temperatures for Main Central Thrust hanging wall
295 mylonites range between ~535°-610°C on the NW Sulej transect, using the Kruhl
296 (1998) opening angle thermometer, and ~ 610° - >735°C based on fabric opening
297 angles and petrologic constraints on the Eastern Sulej transect (Table 1; Law et
298 al., 2013; Stahr, 2013).

299 Dynamic quartz recrystallization microstructures in the Moine Thrust
300 mylonites are similar to microstructures developed under experimental Regime II
301 creep conditions as defined by Hirth and Tullis (1992), with a combination of
302 bulge nucleation recrystallization (BLG) at grain boundaries and subgrain rotation
303 recrystallization (SGR) within quartz grain interiors, and they are consistent with
304 deformation under greenschist facies conditions (Stipp et al., 2002, 2010; Law,
305 2014). Quartz microstructures in the Main Central Thrust mylonites (particularly
306 on the Eastern Sulej transect and at large distances above the thrust on the NW
307 transect) indicate more internal recovery and more extensive grain boundary
308 migration (GBM) than is apparent in the Regime III creep experiments of Hirth
309 and Tullis (1992). The microstructures are more similar to those observed by
310 Stipp et al. (2002; 2010) in quartz veins deformed at natural strain rates and the
311 equivalent of middle-upper amphibolite facies conditions (~550°-700°C; see
312 review by Law, 2014).

313 The small recrystallized quartz grain sizes (16-9 µm), subgrain sizes, and
314 dislocation densities of the Moine Thrust footwall mylonites at the Stack of
315 Glencoul imply large differential stresses (40-250 MPa) during shearing
316 (Weathers et al., 1979; Ord and Christie, 1984); differential stresses of ~ 55-85
317 MPa are recalculated using these grain sizes and the recrystallized grain size



318 piezometer of Stipp et al. (2006) modified by the stress correction of Holyoke and
319 Kronenberg (2010). In contrast, quartz recrystallized grain sizes for Main Central
320 Thrust mylonites on the NW Suttlej transect of 60-35 μm at 75 m above the thrust
321 to 95-75 μm at 750-200 m above the thrust indicate lower differential stresses of
322 19-30 MPa and 13-16 MPa, respectively, using the Stipp et al. (2006) grain size
323 piezometer, adjusted by the same stress correction (Francis, 2012; Law et al.,
324 2013). Flow stresses inferred for the Eastern Suttlej transect mylonites would
325 presumably be even lower than for mylonites of the NW Suttlej transect, but their
326 grain sizes are far greater than those encompassed by any experimental grain size
327 piezometer.

328

329 **3 Methods**

330

331 IR spectroscopy has been an important tool for studying OH defects in nominally
332 anhydrous minerals and, in coordination with deformation experiments, to study
333 the effects of water and hydrogen defects on mechanical properties (e.g., Kats,
334 1962; Griggs and Blacic, 1965; Kekulawala et al., 1978, 1981; Aines and
335 Rossman, 1984; Cordier and Doukhan, 1991; Mackwell and Kohlstedt, 1991; Bai
336 and Kohlstedt, 1996; Kohlstedt et al., 1996). With the introduction of efficient
337 FTIR spectrometers and IR microscopes, studies of intragranular water and
338 hydrogen defects in naturally deformed rocks have been enabled using apertures of
339 50-100 μm (Kronenberg and Wolf, 1990; Nakashima et al., 1995; Gleason and
340 DeSisto, 2008; Finch et al., 2016; Kilian et al., 2016) as well as FTIR mapping of
341 OH contents (Seaman et al., 2013). Molecular water contents of naturally
342 deformed quartz commonly have large grain-to-grain variations at this scale of
343 observation, and it has not been possible to measure OH contents of dynamically
344 recrystallized grains smaller than ~ 75 -130 μm , without including absorption bands
345 of fluid inclusions, fluid films, secondary minerals, and hydrogen species at grain
346 boundaries.



347 Synchrotron-generated IR radiation is much brighter (~1000x) than
348 conventional global sources of commercial FTIR instruments, and high quality IR
349 spectra can readily be measured through small (10 μm) apertures (Lobo et al.,
350 1999; Carr et al., 2008; Ma et al., 2013). By coupling synchrotron radiation to a
351 FTIR instrument and IR microscope, OH absorption bands can be measured
352 through a 10 μm aperture with higher signal-to-noise ratio than through a 100 μm
353 aperture using a standard FTIR. We describe methods of preparing thin
354 unsupported doubly polished IR samples and we compare IR spectra of quartz
355 grains of deformed mylonites collected by FTIR spectroscopy using conventional
356 and synchrotron IR sources. We also describe methods of scanning samples using
357 a 10 μm IR beam, integrating absorbance of selected OH bands (due to fluid
358 inclusions in quartz grains and micas at quartz grain margins and interiors) to form
359 high-resolution images of OH in deformed and recrystallized mylonites. The
360 ability to measure OH bands in finely recrystallized grains is not limited by
361 aperture size or spectral quality for very thin (4–8 μm) samples. Rather,
362 measurements for very thin IR plates are limited by internal reflections that give
363 rise to interference fringes of increasing magnitude with decreasing sample
364 thickness. In practice, measurements of the small OH bands of nominally
365 anhydrous minerals are not possible for IR plate thicknesses of < 25 μm .

366

367 **3.1 Preparation of IR plates**

368

369 Doubly polished IR plates of uniform thickness were prepared perpendicular to
370 foliation and parallel to lineation for Moine Thrust and Main Central Thrust
371 mylonites. Images of large IR plates were collected using a high resolution (4000
372 dpi) Nikon slide scanner (Coolscan 8000 ED), with and without polarizing filters
373 on both sides of the sample. Images of smaller IR plates were recorded with a
374 Zeiss Axioplan 2 petrographic microscope and AxioCam HRc imaging system
375 (Fig. 1). Image contrast was enhanced digitally (using Adobe Photoshop) for



376 crossed-polarized light images of ultrathin samples with small optical retardations.
377 Throughout the preparation of IR plates, impregnation of samples by epoxy or
378 other insoluble resins was avoided to prevent the introduction of OH or CH
379 absorption bands of mounting media, which might be difficult to distinguish from
380 OH bands of samples. While central regions of the IR plates remain fragile, metal
381 frames were mounted to IR plate extremities for mechanical support, to avoid
382 catastrophic loss and facilitate handling.

383 Oriented rock chips were first mounted on a glass thin section plate using
384 CrystalBond 509 (supplied by Aremco, NY) and a low-temperature hot plate. The
385 top surfaces of samples were ground by hand against large glass plates using a
386 sequence of silicon carbide (400 mesh size) and alumina powder slurries of
387 decreasing particle size (600 mesh and 9.5 μm). Grinding with each grit size was
388 carried out until the thickness of the sample, measured by a micrometer, was
389 reduced by three times the grit dimension of the previous step (e.g., 3 μm grit was
390 used to remove $3 \times 9.5 \mu\text{m} = 29 \mu\text{m}$). Samples were cleaned ultrasonically before
391 advancing to the next grit size. Ground surfaces were then polished using
392 polishing cloths (Buehler Texmet and Microcloth) mounted on high-speed laps
393 and a sequence of alumina powder slurries with particle sizes of 3 and 0.3 μm ,
394 using the same measurement technique and testing the finish at each step by visual
395 inspection.

396 Polished specimens were removed from the glass plate and remounted, this
397 time on their polished surfaces, applying gentle pressure to sample centers to
398 promote uniform thickness of the CrystalBond mounting medium ($\sim 20 \mu\text{m}$).
399 Sample plates were cut parallel to the first surface, and the second (cut) surface
400 was ground and polished by the same methods outlined above. For the second
401 surface, a micrometer was used to measure the compound thickness of sample,
402 glass plate, and mounting medium, checking thickness at sample extremities, to
403 adjust the grinding procedure for thickness uniformity. Sample thicknesses were
404 determined from the compound sample-glass plate thickness, the glass plate



405 thickness measured prior to mounting, and estimated CrystalBond layer thickness.
406 In addition to micrometer measurements, sample thicknesses were tested during
407 grinding, using interference colors for thicker samples (first order colors for quartz
408 plates 50-90 μm thick) and maximum detectable birefringence (grey-white scale)
409 while polishing ultrathin ($< 20 \mu\text{m}$) samples.

410 Before removing doubly polished samples from the glass thin section plate,
411 metal frames were mounted to sample extremities using a thin bead of epoxy resin
412 to provide support. For larger ($> 10 \text{ mm}$) samples with thicknesses of 25 - 100
413 μm , metal frames made of Ni wire (1.15 mm diameter) were custom fit to each
414 sample. For thin samples ($< 25 \mu\text{m}$), several metal frames normally used to mount
415 TEM samples (3 mm outer diameter copper rings or slot rings with 1 x 2 mm
416 internal dimensions; supplied by Ted Pella, Inc.) were mounted on each polished
417 sample. Once the metal frames were attached and the epoxy hardened, samples
418 were removed from the glass thin section over a low-temperature hot plate, and the
419 CrystalBond resin removed using acetone.

420 Once quartz IR plates, 4 to 150 μm thick, are removed from the glass thin
421 section plate, they are very fragile, and we found that they did not tolerate
422 ultrasonic agitation. As a result, we soaked samples in acetone baths several
423 times, dissolving CrystalBond resin for 30 minutes for each bath, exchanging the
424 acetone in the beaker and repeating this procedure three times. Even without
425 agitation, some samples did not survive when air bubbles were caught beneath
426 them, leading to specimen warping and disintegration. However, with sufficient
427 care, loss rates were low, and doubly polished IR plates were prepared without
428 impregnating resins with thicknesses of 150 μm to 4 - 8 μm .

429

430 **3.2 FTIR spectroscopy**

431

432 IR absorption spectra were collected for quartz mylonites of varying plate
433 thickness using a Nicolet Magna 560 FTIR with Omnic software, a conventional



434 global source and NicPlan IR microscope (at Texas A&M University), and using a
435 Bruker Hyperion FTIR spectrometer, OPUS software, and IR microscope at the
436 U2A synchrotron beamline of the National Synchrotron Light Source NSLS I (at
437 Brookhaven National Laboratory). Both the Nicolet and Bruker instruments make
438 use of liquid nitrogen-cooled MCT detectors capable of absorbance measurements
439 at wavenumbers of 7400 to 400 cm^{-1} , but we did not routinely store data beyond
440 4000 to 2000 cm^{-1} . We chose to collect spectra at a wavenumber resolution of 4
441 cm^{-1} on both instruments, compiling 512 scans for each spectrum, adjusting to
442 larger numbers of scans with the Nicolet FTIR when needed to reduce noise in the
443 spectrum. For purposes of comparison, all absorbances shown in spectra for
444 quartz grains of varying thickness have been normalized to a common thickness of
445 1 mm.

446 Representative IR spectra of individual quartz grains from Moine Thrust
447 and Main Central Thrust samples measured with the Nicolet FTIR (100 μm
448 aperture) exhibit OH absorption bands of variable size and character (Fig. 2).
449 Moine Thrust quartz spectra show a broad OH stretching band at 3400 cm^{-1} of
450 large but variable magnitude (Fig. 2a), characteristic of molecular water in fluid
451 inclusions of milky quartz (Kekulawala et al., 1978; Aines and Rossman, 1984;
452 Stünitz et al., in press). In addition to this absorption band, some quartz grains
453 from Moine Thrust samples show an additional absorption band (or subtle
454 shoulder) at $\sim 3600 \text{ cm}^{-1}$ (Fig. 2b) due to fine scale micas (sometimes visible
455 optically and sometimes too fine to resolve) dispersed within quartz grains. Main
456 Central Thrust quartz spectra tend to have smaller OH absorption bands (Fig. 2c),
457 with some grains showing small sharp bands at 3595, 3482, 3431, 3408, and 3379
458 cm^{-1} due to hydrogen interstitial defects (Kats, 1962; Aines and Rossman, 1984)
459 and less common grains with a larger broad band at 3400 cm^{-1} . Finely dispersed
460 micas are less common in these coarser-grained mylonites, but some quartz grains
461 also exhibit a 3600 cm^{-1} OH band due to micas (Fig. 2d).



462 Aside from differences in the Nicolet and Bruker FTIR instruments and
463 software, the most significant difference between these facilities is the IR source,
464 so that OH absorption measurements with the Nicolet FTIR and its conventional
465 IR source could not practically be made with apertures < 50-100 μm , while OH
466 absorption bands could routinely be made with the Bruker FTIR and synchrotron
467 IR source through a 10 μm aperture. In both cases, the IR sources are nearly
468 isotropic, leading to differences in those OH absorption bands that are dichroic in
469 quartz and mica grains of varying orientation. Small sharp OH bands due to
470 hydrogen interstitials of quartz between 3595 cm^{-1} and 3379 cm^{-1} absorb most
471 strongly for vibrational directions perpendicular to the quartz *c* axis (Kats, 1962),
472 and these bands vary according to the *c*-axis orientation of each quartz grain
473 relative to the unpolarized IR beam. OH absorption bands at $\sim 3600 \text{ cm}^{-1}$ of micas
474 are strongly dichroic (Beran, 2002), and they vary according to mica grain
475 orientations relative to the unpolarized IR beam. However, the primary OH band
476 of interest for water-weakening of milky quartz, at 3400 cm^{-1} , is broad and
477 isotropic (Kekulawala et al., 1978; Aines et al., 1984), and associated with finely
478 dispersed fluid inclusions. Variations in absorbance of this band for neighboring
479 quartz grains represent real variations in water content and not variations in quartz
480 grain orientation.

481 IR microscope stages are not purged as effectively as sample chambers of
482 FTIR spectrometers that are evacuated or operate at positive nitrogen pressures,
483 and some of our spectra show evidence of atmospheric contamination. We
484 attempted to minimize absorption bands due to H_2O vapor and CO_2 by making
485 spectral measurements of samples and the background (without the sample in the
486 IR beam) at similar times after the last sample change or disturbance at the
487 microscope stage. However, atmospheric absorptions were not always cancelled
488 completely, with the appearance of “noise” between 3940 and 3420 cm^{-1} due H_2O
489 vapor vibrations (Fig. 2a, Grain 26; Fig. 2b, Grains 11B and 39) and absorptions at
490 2337 and 2361 cm^{-1} due to CO_2 , which were either under- or over-corrected by



491 subtracting the measured background (Fig. 2a, Grain 26 and Fig. 2b, Grains 11B
492 and 39, and Fig. 2d, Grain 4 show examples of CO₂ band under-correction; Fig.
493 2c, all grains show negative deflections at 2337 and 2361 cm⁻¹ due to over-
494 correction of CO₂ absorptions).

495 Direct comparisons of intragranular water of a given quartz grain using the
496 two FTIR facilities and a common aperture size could not be made because of 1)
497 the poor signal-to-noise ratio of spectral measurements through a 10 μm aperture
498 with the broad global IR radiation, and 2) the inability to make spectral
499 measurements for a 100 μm-apertured area with the narrow synchrotron IR beam
500 (which is not much broader at the microscope stage than the 10 μm aperture).
501 Samples with uniform OH contents might serve as standards to compare OH
502 absorption bands for different spectrometers, irrespective of aperture size, but our
503 observations indicate that OH contents of the quartz grains in the mylonites we
504 measured are highly variable. Measurements through finer apertures led to larger
505 variances in OH absorption bands, within grains as well as between neighboring
506 grains.

507 We also observed larger amplitude interference fringes in spectra measured
508 with the synchrotron IR source than those observed in spectra measured with the
509 global IR source for a given sample. These fringes are caused by internal
510 reflection within doubly polished, parallel-sided IR plates, and they have larger
511 magnitudes for the highly collimated IR synchrotron beam than for the broad,
512 confocal global IR radiation.

513

514 **3.3 Interference fringes**

515

516 The appearance of interference fringes in spectra can be useful to determine the
517 optical path length in thin samples, as long as their amplitude is small compared
518 with vibrational absorption bands of interest, or they are distinguishable from
519 absorption bands by their wavenumber spacing. Interference fringes are present in



520 many spectra we collected for quartz mylonite plates (Fig. 2c, Grains 1 and 8), that
521 are of manageable size for our thicker samples and problematic for thinner
522 samples.

523 Interference fringes were routinely observed for all samples measured with
524 synchrotron IR radiation and the Bruker FTIR instrument. We therefore made use
525 of these fringes to determine local sample plate thicknesses. For sample plates
526 with large interference fringes, we attempted to reduce their size by tilting samples
527 by 45° in the IR beam or using a Cassegrain objective lens with large numerical
528 apertures (0.6). However, neither of these methods was effective in reducing
529 amplitudes of some of the very large interference fringes sufficiently to detect
530 small OH absorption bands. We had greater success reducing interference fringes
531 by fitting them where the baseline was free of absorption bands, using DatLab
532 software (similar to fringe modeling of Clark and Moffatt, 1978; Pistorius and
533 DeGrip, 2004; Konevskikh et al., 2015), and subtracting the model fringes from
534 the spectral data. This improved the quality of spectra when interference fringes
535 had modest amplitudes, but interference fringes for thin samples were very large
536 and resulting backgrounds sufficiently irregular that we could not resolve small
537 OH absorption bands.

538

539 **3.4 IR plate thickness and OH absorbance determinations**

540

541 IR plate thicknesses were determined by focusing on imperfections in the top and
542 bottom specimen surfaces, and recording the numerical graduation on the focusing
543 knob of the Nicolet FTIR microscope for each surface. The focusing knob
544 graduations were calibrated for vertical translation using the same procedure for a
545 thick (~3 mm) quartz specimen that had been measured with a micrometer.

546 Local IR plate thicknesses were also determined from interference fringes
547 measured in spectra collected with the synchrotron-Bruker FTIR microscope
548 system, where thickness t is given by



549

$$t = 1 / (2 n \delta \nu) \quad (1)$$

551

552 where n is the mean refractive index of quartz ($n = 1.55$) and $\delta \nu$ is the measured
553 peak-to-peak fringe spacing (Stuart et al., 1996; Griffiths and de Haseth, 2007).

554 The two measures of thickness were in agreement within resolution ($\sim 5\%$) for a
555 given IR plate and location within the specimen, with thick IR specimens (~ 100
556 μm) showing real variations in local thickness of $\pm 10 \mu\text{m}$ and thin samples
557 varying in local thickness from 4 to 8 μm .

558 IR spectra were collected and integrated absorbances of OH bands were
559 measured above an assumed straight-line background, where backgrounds were fit
560 and integration limits were chosen at the same wavenumber values, from ~ 3705
561 cm^{-1} to 2880cm^{-1} for OH bands of quartz grains and $\sim 3702 \text{cm}^{-1}$ to 3544cm^{-1} for
562 OH bands of micas (Fig. 3a). We have confidence in our ability to separate OH
563 absorption bands due to fluid inclusions in quartz and due to micas in Moine
564 Thrust samples, because IR spectra of coarse muscovite grains in Moine Thrust
565 mylonites consist of a single OH absorption band at 3620cm^{-1} (Fig. 3b). This
566 strong OH stretching band is well known from multiple spectroscopic studies of
567 muscovite (Beran, 2002; Tokiwai and Nakashima, 2010a,b; Kallai and Lapides,
568 2015).

569 Our ability to distinguish OH absorbance due to fluid inclusions and micas
570 in Main Central Thrust samples, by contrast, is poor. In addition to the primary
571 OH absorption bands of muscovite (Fig. 3c) and biotite (Fig. 3d) grains at 3638
572 cm^{-1} and 3614cm^{-1} , respectively, Main Central Thrust muscovite spectra show
573 smaller OH bands at 3815 , 3311 , 3146 , and 3035cm^{-1} and biotite spectra show
574 shoulders at both sides of the primary OH absorption band (at $\sim 3679 \text{cm}^{-1}$ and
575 $\sim 3561 \text{cm}^{-1}$) and significant OH bands at 3258 , 3043 , and 2829cm^{-1} . All of these
576 OH bands are anisotropic, and complexly so. OH band absorbances of micas will
577 therefore vary with crystallographic orientation within an unpolarized IR beam.



578 The primary OH band of Main Central Thrust muscovite at 3638 cm^{-1} , measured
579 in polarized IR radiation, is strongest when the vibration direction E is parallel to
580 the basal plane (001), consistent with idealized hydrogen positions in dioctahedral
581 micas and the polarization of OH bands normally reported for muscovite (Beran,
582 2002). However, our polarized IR measurements of muscovite OH bands at 3311,
583 3146, and 3035 cm^{-1} indicate that they are more nearly isotropic. The primary OH
584 band of Main Central Thrust biotite grains at 3614 cm^{-1} is strongest when E is
585 perpendicular to (001), consistent with hydrogen positions of trioctahedral micas
586 and OH band polarizations observed for phlogopite (Beran, 2002). However, the
587 OH bands at 3258, 3043, and 2829 cm^{-1} are only weakly polarized and in the
588 opposite sense of the primary 3614 cm^{-1} band.

589 As a result, we only feel confident in our determinations of OH
590 absorbances of quartz grains in Main Central Thrust mylonites (Fig. 2c), when
591 mica inclusions are absent (3600 cm^{-1} band is undetectable). When the 3600 cm^{-1}
592 band is present (Fig. 2d), we cannot readily interpret absorption bands of
593 molecular water or hydrogen defects in quartz; at best, quartz OH contents are
594 overestimated.

595 Integrated absorbances of OH bands of quartz were used to determine OH
596 contents based on the Beer-Lambert relation (Stuart et al., 1996; Griffiths and
597 Haseeth, 2007),

598

$$599 \quad A = k c t \quad (2)$$

600

601 where integrated absorbance Δ ($= A/t$, determined in cm^{-2}), is related to the
602 concentration c of OH, assuming that k for the broad isotropic OH band at 3400
603 cm^{-1} due to fluid inclusions in milky quartz is the same as for the broad isotropic
604 OH absorption band of molecular water in wet synthetic quartz (Aines et al.,
605 1984);

606



607 c (in molar ppm, OH/10⁶ Si) = 1.05 Δ (in cm⁻²). (3)

608

609 To the extent that spectra include sharp OH absorption bands due to hydrogen
610 point defects, this calibration will overestimate OH concentrations due to H
611 interstitials, given that the value of k for these bands is larger than for molecular
612 water (Kats, 1962; Thomas et al., 2009) for quartz grains oriented for maximum
613 OH absorbance (vibrational directions perpendicular to the c axis). However,
614 apparent OH concentrations due to H interstitials will appear smaller for quartz
615 grains oriented for minimum OH absorbance in unpolarized IR radiation. These
616 errors are not serious for quartz grains with large molecular water contents that
617 dominate over hydrogen point defect concentrations. However, we expect that our
618 determinations of OH concentrations for dry quartz grains (< 100 ppm) with
619 spectra dominated by sharp dichroic OH bands are not as accurate as for wet
620 quartz grains with spectra dominated by isotropic broad OH bands. We also
621 acknowledge that the calibration for molecular water absorption of synthetic
622 quartz used here differs from the calibration of the 3400 cm⁻¹ OH band absorbance
623 reported by Stipp et al. (2006) for milky quartz grains with fluid inclusions. If this
624 alternative calibration is correct, OH contents reported here are smaller by a factor
625 of 0.56. The calibration of Stipp et al. (2006) is based on FTIR measurements of
626 individual grains of a pure quartzite and independent Karl Fischer titration of
627 disaggregated quartz particles, and we do not understand the source of
628 discrepancies in absorbance calibrations.

629 Integrated absorbances Δ (A/t in cm⁻²) were also measured for the 3600 cm⁻¹
630 OH band of micas but no attempt was made to convert these to OH
631 concentrations. Muscovite grains of Moine Thrust samples show simple OH
632 absorption spectra with one prominent band at ~3600 cm⁻¹ (Fig. 3b), but integrated
633 absorbances for this band cannot be converted to muscovite (or OH) content
634 without polarized IR spectral measurements and detailed knowledge of muscovite
635 grain orientations. Integrated absorbances of OH bands of muscovite and biotite



636 grains in Main Central Thrust samples at wavenumbers of 3800 cm^{-1} to 2830 cm^{-1}
637 are more complicated to interpret. In addition to the need for polarized IR
638 measurements and grain orientation information, integrated OH absorbances due
639 to unseen (below optical resolution) mica inclusions will depend on whether the
640 micas consist of muscovite or biotite, or both.

641

642 **3.5 Mapping OH absorption bands**

643

644 Integrated areas of OH were measured as a function of spatial x and y dimensions
645 within IR plates measured on the Bruker FTIR system, separating absorbances of
646 OH bands due to fluid and mica inclusions to the extent that this was possible.
647 Samples were translated under the $10\text{ }\mu\text{m}$ IR beam using a motorized x-y stage,
648 controlled by OPUS software, and spectra were collected for each translation step,
649 usually moving the stage by $10\text{ }\mu\text{m}$. In those cases where we mapped coarse-
650 grained samples over large regions (up to $\sim 4.5\text{ mm}^2$), we increased the step
651 spacing to 30 or $50\text{ }\mu\text{m}$, in order to limit total step-scan measurement times and the
652 number of required steps (our software version could not compile > 99 steps in a
653 single row of measurements). Integrated absorbance measurements were made
654 over $3705 - 2880\text{ cm}^{-1}$ (to include the 3400 cm^{-1} broad band of milky quartz) and
655 $3702 - 3544\text{ cm}^{-1}$ (to determine the 3600 cm^{-1} mica band) for each scanned area
656 (compiling absorbance measurements for 660 to 4950 spectra). The integrated
657 absorbances of the 3600 cm^{-1} mica band were subtracted from the first of the two
658 integrated absorbances, saving the result as a representative measure of the OH
659 bands of molecular water and hydrogen defects of quartz.

660 Given that the same IR background was used to reference all spectral
661 measurements made during x-y mapping of OH bands, long step-scan mapping
662 projects were susceptible to changes in aperture and condenser lens centering, IR
663 beam drift and beam outages while the synchrotron ring was refilled. As a result,
664 we modified our IR absorbance measurement methods for step-scan mapping to



665 obtain reasonably accurate integrated absorbances of broad OH bands without
666 requiring high resolution or precision measurements of small, sharp OH bands,
667 with the goal of limiting total measurement times.

668 IR spectra used to map OH bands were made at a wavenumber resolution
669 of 8 cm^{-1} , reducing the interferometer mirror translation times with no detectable
670 degradation in integrated absorbance measurements. More significantly, we
671 reduced the number of scans to measure each IR spectrum from 512 to just 16,
672 resulting in measurement times for each spectrum of only $\sim 10\text{ s}$, including the
673 time for stage translation. This number of scans was chosen on the basis of tests
674 for a given quartz grain of Moine Thrust sample SG-10 with an absorbance of 36
675 cm^{-1} measured over $3705 - 2880\text{ cm}^{-1}$ (or $\Delta = 3000\text{ cm}^{-2}$ for a sample thickness t of
676 $120\text{ }\mu\text{m}$), determining integrated absorbances for spectra compiled for 512, 256,
677 128, 64, 32, 16, and 8 scans. Integrated absorbances of the broad 3400 cm^{-1} band
678 varied by less than 4% for measurements made with 512 - 32 scans, while the
679 absorbance measurements for 16 and 8 scans were within 8% and 12%,
680 respectively, of the mean of the other determinations. OH maps were thereby
681 constructed for sample areas of 0.066 mm^2 ($22 \times 30 = 660$ spectra) to nearly 4.5
682 mm^2 ($50 \times 99 = 4950$ spectra) requiring measurement times of 110 minutes to
683 nearly 14 hours, respectively. The larger step-scan maps include some poor
684 spectral data corresponding to beam drift or interruption of radiation; these
685 spectra were identified and removed from the data set.

686 Maps of integrated OH absorbance were contoured (using SigmaPlot) to
687 form images of water contents and micas, and superposed on optical micrographs
688 of the measured regions, to look for correlations between OH content and
689 microstructure. OH contents of quartz, due primarily to the broad 3400 cm^{-1} OH
690 band of fluid inclusions and secondarily to sharp hydrogen defect bands, were
691 mapped as molar ppm ($\text{OH}/10^6\text{Si}$), while the OH absorption band at $\sim 3600\text{ cm}^{-1}$
692 due to micas was mapped in units of integrated absorbance (cm^{-2}). Logarithmic
693 contour intervals of $\log_{10}(\text{OH}/10^6\text{Si})$ were chosen to image water contents in quartz



694 over a wide range of concentrations, and to provide visual images of water
695 distributions at low (and high) values. OH contents were contoured in color with
696 blue (and other cool colors) corresponding to large water contents, and red (and
697 other warm colors) corresponding to low water contents. Logarithmic contours
698 were also chosen to image distributions of micas, plotting $\log_{10}(\Delta \text{ in cm}^{-2})$ for the
699 integrated absorbance of the $\sim 3600 \text{ cm}^{-1}$ OH band, using a similar key for the
700 contour interval, where cool colors correspond qualitatively to high mica contents
701 and warm colors correspond to low mica contents.

702

703 **4 Results**

704

705 Synchrotron radiation coupled with the FTIR spectrometer and microscope at the
706 NSLS I has enabled spectral measurements for sampling volumes that are smaller
707 by a factor of 100 than is possible with a conventional FTIR microscope system.
708 OH band measurements for quartz through a $10 \mu\text{m}$ aperture are comparable to
709 OH band measurements made using a conventional FTIR spectrometer-
710 microscope system through a $100 \mu\text{m}$ aperture. The ability to measure OH bands
711 of individual recrystallized quartz grains is currently limited by IR plate thickness,
712 not because of signal-to-noise problems, but because of internal reflections in
713 ultrathin IR plates. Water contents of quartz grains, and OH contents associated
714 with micas and other hydrous phases, can be measured spatially by translating
715 samples under the $10 \mu\text{m}$ -apertured beam and contouring integrated OH
716 absorbance over chosen wavenumber intervals. High-resolution images, formed
717 in this way, can be compared with optical deformation microstructures, providing
718 a means of studying the role of water in deformation and changes in water content
719 during deformation and recrystallization.

720

721 **4.1 Synchrotron FTIR Measurements**

722



723 IR spectra of quartz in Moine Thrust and Main Central Thrust samples measured
724 with the synchrotron-FTIR system through a 10 μm aperture exhibit OH
725 absorption bands of similar character at the same wavenumbers as measured for
726 samples of the same thickness ($\sim 100 \mu\text{m}$) using a conventional FTIR
727 spectrometer-microscope system through a 100 μm aperture. OH bands of Moine
728 Thrust quartz grains generally include a large broad absorption band at 3400 cm^{-1}
729 due to dispersed fluid inclusions, with some grains showing a secondary
730 absorption band at 3600 cm^{-1} due to mica inclusions (Fig. 4a), just as those
731 measured for larger sampling volumes (Fig. 2a, b). Given the more intense IR
732 radiation of the synchrotron source, we did not encounter any losses in spectral
733 quality due to signal-to-noise ratios for the smaller sample volumes measured.
734 However, interference fringes were more apparent using the synchrotron-FTIR
735 microscope system associated with internal reflections from the parallel-sided IR
736 plates and the lower divergence of the synchrotron beam. Synchrotron FTIR
737 measurements of Main Central Thrust samples through a 10 μm aperture show the
738 same sharp OH bands due to hydrogen point defects as measured through a 100
739 μm aperture. With the exception of some unusual grains, broad OH bands at 3400
740 cm^{-1} in Main Central Thrust samples are small to negligible.

741 Integrated absorbances of the broad 3400 cm^{-1} OH band measured with a
742 10 μm aperture are highly variable spatially, yielding water contents for individual
743 spots within quartz grains for a given Moine Thrust sample (SG-10) of 280 to
744 9000 ppm, with a comparable mean (2430 ppm) but more variance among
745 individual measurements than measured with a 100 μm aperture (from 1130 to
746 8590 ppm for 45 measurements). For dry Main Central Thrust samples (S09-
747 71B), OH contents of quartz measured using a small (10 μm) aperture varied from
748 50 to 300 ppm while OH contents measured with the larger (100 μm) aperture
749 varied from 85 to 240 ppm. For unusually wet Main Central Thrust samples (S09-
750 35), OH contents of quartz measured using a small (10 μm) aperture varied from
751 150 to 7500 ppm while OH contents measured with the larger (100 μm) aperture



752 varied from 160 to 4620 ppm. The variations in size of sharp OH bands associated
753 with hydrogen point defects, by comparison, are much smaller. The sizes of these
754 dichroic absorption bands may be explained by differing quartz grain orientations
755 without calling upon any variations in hydrogen point defect concentrations
756 (Thomas et al., 2009).

757

758 **4.2 Limits of IR plate thickness**

759

760 IR spectra of quartz were measured through a 10 μm aperture, varying sample
761 plate thickness from $\sim 100 \mu\text{m}$ to 4–8 μm , with the hope that we might be able to
762 measure the OH bands of individual recrystallized grains in ultrathin IR plates.
763 Owing to the coarse recrystallized grain sizes of Main Central Thrust mylonites,
764 we were able to measure IR spectra of most quartz grains in these samples for
765 plate thicknesses of $>50 \mu\text{m}$. However, our IR measurements of individual grains
766 in Moine Thrust samples (with recrystallized grains of $\sim 10\text{--}15 \mu\text{m}$) are limited to
767 larger porphyroclastic and deformed ribbon quartz grains. While the signal-to-
768 noise ratio of our measurements continue to be acceptable to measure the small
769 OH absorption bands in very thin IR plates, interference fringes increase in size as
770 the IR plate thickness t is decreased. Interference fringes become very large at $t <$
771 $25 \mu\text{m}$ and OH bands cannot be detected at $t = 13 \mu\text{m}$ and $6.5 \mu\text{m}$ (Fig. 4b, c).

772 The magnitude of interference fringes can be modified by tilting the
773 sample within the IR beam and increasing the numerical aperture of the
774 Cassegrain objective on the IR microscope. However, with peak-to-trough fringe
775 magnitudes 10^4 greater than the OH bands in our mylonite samples, the modest
776 reductions in fringe amplitude realized by these methods are not significant for the
777 measurement of OH absorptions. Our efforts to model interference fringes were
778 largely successful, reducing their magnitudes by a factor of ~ 100 , but the resulting
779 backgrounds were not flat enough at the scale of the known OH absorption bands
780 of the Moine Thrust (or Main Central Thrust) samples to allow absorption band



781 measurements for our thinnest sample plates. Interference fringes vary in
782 magnitude according to IR plate thickness, surface polish, and scattering by
783 internal flaws, but OH absorption bands could only be measured for samples of
784 ~40 μm thickness or greater.

785

786 **4.3 Optical microstructures and plate thickness**

787

788 Although the painstaking efforts to prepare ultrathin IR plates were not rewarded
789 by spectral measurements of OH content within individual, finely recrystallized
790 grains, optical imaging of deformation and recrystallization microstructures in
791 ultrathin samples was improved over that using conventional thin sections.
792 Optical microstructures of 100 μm -thick IR plates are poorly resolved by
793 comparison with those imaged in 30 μm -thick sections (Fig. 1c, d), with
794 interference colors that reflect greater optical retardation, grain boundaries that are
795 not as clearly defined, and greater numbers of overlapping grains. Remarkably,
796 ultrathin IR plates of quartz mylonites, just 4-8 μm thick (Fig. 1f), continue to
797 exhibit contrast between grains and within grains, with first-order black to grey
798 birefringence that can be enhanced by increasing image contrast.

799 High magnification optical micrographs of ultrathin Moine Thrust
800 mylonites reveal microstructures (Fig. 5) that correspond better to TEM
801 observations (Weathers et al., 1979; Ord and Christie, 1984) than to optical
802 microstructures of conventional thin sections. While highly flattened ribbon
803 quartz grains observed in 30 μm sections show smoothly varying undulatory
804 extinction, ultrathin sections exhibit well developed subgrains within grain
805 interiors, with sharply defined changes in extinction marking the locations of
806 distinct subgrain walls (Fig. 5a). Prior TEM of Moine Thrust mylonites revealed
807 significant densities of free dislocations that have been associated with smooth
808 changes in extinction. However, TEM observations also show dense, sharply
809 defined subgrain walls (Weathers et al., 1979; Ord and Christie, 1984) that are



810 often not resolved as optical microstructures of normal thin sections. These
811 observations suggest that ultrathin sections may be useful to distinguish smooth
812 changes in extinction (due to free dislocations and strains internal to subgrains)
813 from discrete changes in extinction of overlapping and neighboring subgrains.
814 Recrystallized grains are apparent at some sutured grain boundaries in
815 Moine Thrust samples (Fig. 5a) with grain sizes that are significantly smaller than
816 the dimensions of subgrains within the ribbon quartz grains. Microstructures of
817 these new grains suggest that they form by grain boundary bulge mechanisms.
818 Newly recrystallized grains surrounding other grains (Fig. 5b) have sizes in
819 common with nearby subgrains, and microstructures suggest that these
820 recrystallized grains develop by subgrain rotation. While prior optical
821 examination of Moine Thrust mylonites using 30 μm thin sections has led to the
822 conclusion that both BLG and SGR recrystallization were important during
823 deformation (Law, 2014), quantitative evaluations of these processes and their
824 contributions to dislocation creep could be improved by the higher microstructural
825 resolution offered by ultrathin sections.

826

827 **4.4 High-resolution imaging of OH**

828

829 By translating IR plates under the 10 μm apertured IR beam and measuring spectra
830 over many steps (spaced by 10 to 50 μm) we were able to compile integrated
831 absorbances over spatial areas of up to 4.455 mm^2 (50 x 99 = 4950 spectra, step
832 size = 30 μm). Integrated absorbances of OH bands were determined for the
833 collection of spectra, choosing limits of 3705 cm^{-1} to 2880 cm^{-1} to include the
834 broad OH band at 3400 cm^{-1} and sharp OH bands between 3595 cm^{-1} and 3379 cm^{-1}
835 ¹ due to fluid inclusions and hydrogen defects, respectively, within quartz grains.
836 Integrated absorbances of the 3600 cm^{-1} OH band of micas were determined
837 between 3702 cm^{-1} to 3544 cm^{-1} , and this value, $\Delta_{\text{OH musc}}$ was subtracted from the



838 first integrated absorbance (as illustrated in Fig. 3a) to represent the best measure
839 of absorbance $\Delta_{\text{OH}_{\text{qtz}}}$ of water and hydrogen defects of quartz grains.

840 These values were plotted spatially for Moine Thrust and Main Central
841 Thrust samples, and contoured on common logarithm scales, to form high-
842 resolution images (Figs. 6 - 9) of OH absorbance of quartz (converting to molar
843 ppm, $\text{OH}/10^6\text{Si}$) and OH absorbance of micas ($\Delta_{\text{OH}_{\text{musc}}}$ in cm^{-2}). In all cases, we
844 superposed the contoured OH maps on optical micrographs to make sense of OH
845 distributions in terms of fluid inclusions, defects, and micas, and their
846 relationships to deformation and recrystallization microstructures. Images of the
847 3600 cm^{-1} mica OH absorbance were placed over plane light micrographs that
848 highlight scattering (and sometimes color) due to micas, grain boundaries, and
849 fluid inclusions, while images of OH absorbance due to fluid inclusions
850 (principally the broad 3400 cm^{-1} band at large integrated absorbances) and lesser
851 hydrogen defects of quartz were placed over cross-polarized light micrographs that
852 emphasize undulatory extinction, deformation microstructures and
853 recrystallization of quartz grains.

854 FTIR maps of OH of quartz grains show that water contents of deformed
855 mylonites (Figs 6 - 8) are extremely heterogeneous when absorbance is dominated
856 by the 3400 cm^{-1} band, with water contents that vary from 300 ppm to $>10,000$
857 ppm between neighboring grains and within grain interiors. This result helps
858 explain the wide variations in water contents measured by conventional FTIR
859 spectroscopy through a $100\text{ }\mu\text{m}$ aperture. OH contents of quartz grains of
860 relatively dry mylonites (Fig. 9) are more nearly constant (50 - 150 ppm) when IR
861 spectra are dominated by sharp OH bands of hydrogen point defects. FTIR maps
862 of mylonite samples, constructed for the 3600 cm^{-1} OH absorbance, show that
863 micas are heterogeneously distributed, with concentrations apparent at quartz grain
864 boundaries in optical micrographs (Figs. 6 - 9) and finely dispersed micas within
865 quartz grain interiors (Figs. 7, 8), some of which can be difficult to detect



866 optically. We show these relationships in several contoured OH maps, focusing
867 on different types of quartz grains and deformation microstructures.

868 While most quartz grains in Moine Thrust mylonites are highly deformed,
869 some quartz grains remained nearly equant (globular grains or quartz augen with
870 *c*-axes aligned perpendicular to foliation; Law et al., 1986), despite their high
871 original OH contents in the form of fluid inclusions (Fig. 6). The lack of 3600 cm^{-1}
872 ^1OH absorbances within globular quartz grains indicate that finely dispersed
873 micas in original relict quartz grains are absent, while large 3600 cm^{-1} OH
874 absorbances at quartz grain extremities indicate that micas are localized at quartz
875 grain boundaries (Fig. 6a, b). Broad band OH contents of quartz of >1000 ppm,
876 thought to be sufficient for water weakening, are present in undeformed and
877 deformed ribbon quartz grains (Fig. 6 c, d) with very large water contents
878 ($>10,000$ ppm) marking a healed crack in the quartz augen, made up of a planar
879 array of fluid inclusions. OH contents due to fluid inclusions are also very large at
880 the globular augen quartz boundaries, coincident with high mica concentrations,
881 and in some ribbon and recrystallized quartz grains. With a sample plate thickness
882 of $56\text{ }\mu\text{m}$, these maps represent OH absorbances within the interiors of larger
883 quartz grains, while OH absorbances of fine micas and recrystallized quartz
884 represent composite spectra of polycrystalline fault rock. However, we are
885 confident that the quartz OH contents reflect fluid inclusions, even in these fine-
886 grained regions, because of the simple spectral quality of muscovite (Fig. 3) in the
887 Moine Thrust samples and our ability to distinguish between the 3400 cm^{-1} and
888 3600 cm^{-1} OH absorbances.

889 High-strain ribbon quartz grains in the Moine Thrust samples have large
890 OH contents (>1000 ppm) comparable to those of water-weakened synthetic and
891 milky quartz (samples SG-10 and SG-8, Figs. 7 and 8, respectively), with some
892 reductions in OH at recrystallized margins of original grains. Micas, as imaged by
893 the 3600 cm^{-1} OH absorbance (Fig. 7a, b), continue to be highly localized at the
894 grain boundaries of deformed quartz grains, with a mixture of fine-grained mica



895 and quartz grains providing evidence for some redistribution during
896 recrystallization (Fig. 8a, b). The broad 3400 cm^{-1} OH absorbance in marginal
897 recrystallized regions surrounding ribbon quartz grains are locally smaller than
898 those of the original deformed quartz grains in some regions (Fig. 7 c, d) while
899 broad OH absorbances of recrystallized quartz continue to be large where mica
900 contents (as evidenced by the 3600 cm^{-1} band absorbance) are large (Fig. 8 c, d).

901 FTIR maps of coarse-grained Main Central Thrust mylonites yield spectral
902 measurements of individual grains of quartz, muscovite and biotite, even for
903 relatively thick IR plates and larger step sizes (Fig. 9). Small, sharp absorption
904 bands of quartz grains yield OH contents of ~ 100 ppm, with only local regions of
905 quartz with larger OH contents near contacts with coarse-grained muscovite and
906 biotite grains. Quartz grain interiors generally lack the absorbance band at 3600
907 cm^{-1} ; thus, there is no evidence for finely dispersed micas within these coarse-
908 grained deformed (and recrystallized) quartz grains. Contours of integrated OH
909 absorptions at 3400 cm^{-1} are considerably larger for coarse-grained micas and near
910 their contacts with quartz grains. However, these bands cannot be attributed to
911 fluid inclusions where they coincide with large 3600 cm^{-1} mica bands, given that
912 muscovite and biotite grains in the Main Central Thrust samples exhibit complex
913 secondary OH bands between 3311 and 2829 cm^{-1} .

914

915 **5 Discussion**

916

917 Our IR spectra collected from quartz-rich mylonites in the footwall to the Moine
918 Thrust (Scottish Caledonides) and the hanging wall to the Main Central Thrust
919 (Himalaya of NW India) using synchrotron IR radiation through a $10\text{ }\mu\text{m}$ aperture
920 are comparable to IR spectra we collected for the same samples using a
921 conventional FTIR-microscope system through a $100\text{ }\mu\text{m}$ aperture. The broad OH
922 band and large water contents of the Moine Thrust mylonites deformed at 300 -
923 350°C are in line with previous FTIR studies of OH in quartz deformed under



924 greenschist facies conditions. The sharp OH bands and low water contents of the
925 Main Central Thrust mylonites deformed at 570 - 610°C are consistent with FTIR
926 studies of quartz in other shear zones deformed at amphibolite conditions (Han et
927 al., 2013; Kilian et al., 2016).

928 This study adds to an emerging impression that quartz tectonites deformed
929 at greenschist temperatures and natural strain rates are wet. Water contents of
930 quartz mylonites from other locations of the Moine Thrust deformed at greenschist
931 facies conditions vary from 1400 to 7500 ppm (Kronenberg and Wolf, 1990).
932 Water contents of deformed quartz across the Median Tectonic Line (Japan) vary
933 from 300 to 2500 ppm (Nakashima et al., 1995), depending on metamorphic grade
934 and shear displacement. Water contents of quartz deformed in granitic shear zones
935 and mylonites at greenschist conditions reach values of 1100 ppm and >10,000
936 ppm (Gleason and DeSisto, 2008; Kronenberg et al., 1990).

937 Our measurements of water in quartz grains in Main Central Thrust
938 mylonite samples build on the results of Nakashima et al. (1995), Han et al. (2013)
939 and Kilian et al. (2016), which suggest that quartz deformed at amphibolite
940 temperatures and natural strain rates may be relatively dry. Han et al. (2013)
941 report water contents of just 10 to 110 ppm for quartz grains of granitic mylonites
942 deformed at 400-500°C within the Longmenshan tectonic zone (Sichuan, China).
943 Kilian et al. (2016) measured OH contents of just 20 to 100 ppm for quartz grains
944 of granitic mylonites deformed at upper amphibolite conditions (Truzzo meta-
945 granite, Central Alps, Italy), with IR spectra dominated by sharp OH bands due to
946 hydrogen point defects and no detectable broad band due to molecular water.
947 Significant water contents (1400-4400 ppm) have been observed for quartz
948 deformed at amphibolite conditions in mylonites of the El Pichao shear zone (NW
949 Argentina), though reductions in water content are evident with progressive
950 deformation (Finch et al., 2016). Fluid inclusions may be lost and intragranular
951 water contents reduced during high temperature deformation by a variety of
952 processes, from pipe diffusion (Bakker and Jansen, 1990, 1994; Hollister, 1990;



953 Cordier et al., 1994; Mavrogenes and Bodnar, 1994) to recrystallization and grain
954 boundary sweeping (Faleiros et al., 2010; Seaman et al., 2013; Finch et al., 2016;
955 Kilian et al., 2016), and partial melting (Seaman et al., 2013).

956

957 **5.1 Wide variations in OH content**

958

959 All previous FTIR studies of quartz mylonites have revealed large variations in
960 water content for different grains of the same fault rock, and FTIR mapping of OH
961 (Seamon et al., 2013, and this study) has revealed significant variations within
962 grains. Given that the broad OH band of quartz at 3400 cm^{-1} is isotropic, the
963 grain-to-grain variations (at large water contents) are real, not owing to variations
964 of dichroic OH band absorbance with quartz grain orientation for a uniform
965 concentration of structurally incorporated hydrogen defects. Our FTIR
966 measurements of quartz OH contents using a $10\text{ }\mu\text{m}$ aperture show that these
967 variations depend on sampling volume, as observed by Kilian et al. (2016) who
968 showed that broad OH band absorptions scale with size of the measurement area
969 due to the inhomogeneous distributions of fluid inclusions.

970 The large and variable water contents of quartz mylonites are far above
971 equilibrium solubilities (e.g., Paterson, 1986; Kronenberg et al., 1986; Cordier and
972 Doukhan, 1989), and the variations in non-equilibrium OH contents probably
973 reflect some part of the history of water migration during deformation. Images of
974 contoured OH absorbances of the broad 3400 cm^{-1} band for quartz and the 3600
975 cm^{-1} bands of micas constructed for Moine Thrust and Main Central Thrust
976 samples show relationships with optical deformation and recovery microstructures
977 that suggest mechanisms by which water is incorporated in quartz grain interiors,
978 how water becomes redistributed during deformation and recovery, and how water
979 is lost from quartz grain interiors.

980 OH contours within relatively undeformed quartz augen of the Moine
981 Thrust reveal planar zones of high water content that correspond to secondary



982 fluid inclusions at healed microcracks (Fig. 6). We know little of this early brittle
983 deformation, but these fluid inclusion arrays have microstructures similar to those
984 generated during hydrothermal diffusional healing of cracks (Smith and Evans,
985 1984; Beeler and Hickman, 2015). As a result, early brittle deformation,
986 infiltration of water along open cracks, and crack healing appears to be important
987 to the early introduction of water to quartz grain interiors (Kronenberg et al., 1986,
988 1990, 2001; Diamond et al., 2010; Tarantola et al., 2010, 2012; Stünitz et al., in
989 press).

990 OH contours in plastically deformed regions of Moine Thrust samples are
991 complex and water contents appear to vary with strain patterns and
992 recrystallization (Figs. 6 - 8). Clearly defined planar arrays of coarse fluid
993 inclusions are absent from these regions, although the measured high water
994 contents indicate redistribution of water within quartz grains rather than the loss of
995 water. Processes of fluid inclusion decrepitation under deviatoric stresses have
996 been studied experimentally (Diamond et al., 2010; Tarantola et al., 2010, 2012,
997 Stünitz et al., in press), and they include simultaneous shrinkage of coarse (> 10
998 μm , optical-scale) inclusions, generation of dislocations at fluid inclusion walls,
999 and formation of a new population of very fine (< 100 nm) fluid inclusions (visible
1000 only by transmission electron microscopy), which reside at dislocations and
1001 resemble water clusters and inclusions of deformed and heat-treated synthetic
1002 quartz (McLaren and Hobbs, 1972; White, 1973; Kirby and McCormick, 1979;
1003 Christie and Ord, 1980). The loss of coarse fluid inclusions and growth of fine
1004 inclusions requires diffusive transport, which may occur along interconnected and
1005 mobile dislocations (McLaren et al., 1983; 1989; Cordier et al., 1988, 1994;
1006 Bakker and Jansen, 1990, 1994; Hollister, 1990; Kronenberg et al., 1990, 2001;
1007 Tarantola et al., 2010, 2012; Stünitz et al., in press). Once formed, very fine fluid
1008 inclusions may also coarsen by pipe diffusion and processes documented by
1009 McLaren et al. (1983) and Cordier et al. (1988), leading to continuous changes in
1010 inclusion densities and size distributions.



1011 Some regions of recrystallized quartz grains in Moine Thrust samples
1012 appear to have somewhat lower OH contents than original ribbon quartz grains
1013 (Figs. 7), while highly recovered and recrystallized quartz grains in the Main
1014 Central Thrust samples have little or no detectable molecular water (Fig. 9).
1015 Reductions in intragranular water during dynamic recrystallization have been
1016 attributed to sweeping of fluid inclusions by mobile grain boundaries and losses of
1017 water from the fault rock by rapid grain boundary diffusion (Faleiros et al., 2010;
1018 Seaman et al., 2013; Finch et al., 2016; Kilian et al., 2016).

1019

1020 **5.2 Limits on measurement volume**

1021

1022 While synchrotron IR radiation has enabled spectral measurements through small
1023 (10x10 μm) apertured areas, we have not succeeded in measuring OH absorption
1024 bands for ultrathin samples when their interference fringes are larger than the
1025 absorption bands. Interference fringes due to internal reflections can be modeled
1026 (e.g., Clark and Moffatt, 1978; Pistorius and DeGrip, 2004; Konevskikh et al.,
1027 2015), but they become increasingly large as thickness of the doubly polished
1028 sample is reduced. For samples < 25 μm thick, we were unable to execute fringe
1029 corrections with sufficient confidence to measure OH spectra of individual,
1030 recrystallized quartz grains of the Moine Thrust.

1031 The physical means by which we attempted to reduce interference fringes
1032 (rotating the sample within the IR beam, and changing numerical aperture of the
1033 IR objective) were also unsuccessful. However, interference fringes might be
1034 reduced in future studies by a number of other methods. Interference fringes could
1035 be eliminated if ultrathin samples are mounted on a substrate with a matching
1036 refractive index n (*personal communication*, Sina Marti). Such a substrate would
1037 need to be infrared-transparent, facilitate bonding between the sample and
1038 substrate (by low temperature melting or casting, to eliminate any residual air gaps
1039 between the sample and substrate), and lack OH bonds of its own. Alternatively,



1040 internal reflections could be reduced if thin samples are immersed in
1041 polychlorotrifluoroethylene oil ($n = 1.41$), which exhibits strong absorption bands
1042 at $<2500\text{ cm}^{-1}$ but has only small bands (between 3600 and 3200 cm^{-1}) that might
1043 interfere with OH absorptions of the sample (*personal communication*, Jed
1044 Mosenfelder and George Rossman). Internal reflections could also be reduced if
1045 only one surface of the ultrathin IR plate is polished, leaving the other surface
1046 precision-ground for a given plate thickness t (as employed by Woodhead et al.,
1047 1991). Spectra measured with just one polished surface will suffer some signal
1048 loss and spectra will need to be corrected for background. Another method of
1049 reducing interference fringes that obscure OH absorption bands might be
1050 developed by focused ion beam (FIB) milling of one side of a doubly polished
1051 sample to eliminate reflections at that surface over the IR wavenumbers of interest
1052 (*personal communication*, Roy Christoffersen).

1053

1054 **5.3 Water weakening in nature?**

1055

1056 The water contents of Moine Thrust samples and other quartz mylonites deformed
1057 at greenschist conditions are large and variable, comparable to (and even larger
1058 than) those of wet varieties of synthetic and natural milky quartz (350-4000 ppm)
1059 that exhibit water weakening in laboratory studies (e.g., Griggs and Blacic, 1965;
1060 Kekulawala et al., 1978; Stünitz et al., in press). Highly deformed ribbon quartz
1061 grains, less deformed quartz augen, and recrystallized quartz grains of the Moine
1062 Thrust exhibit OH bands of similar character to those of quartzites deformed in
1063 laboratory experiments, a result that validates applications of wet quartzite
1064 rheologies to evaluate rates of dislocation creep in middle to upper crustal shear
1065 zones (e.g., Hirth et al., 2001; Behr and Platt 2011; Law, 2014).

1066 In contrast, water contents of highly sheared and recrystallized quartz in
1067 the hanging wall of Main Central Thrust samples and other mylonites deformed at
1068 upper amphibolite conditions are much smaller than measured for wet varieties of



1069 quartz. The large shear strains accumulated over ~150-250 km of displacement
1070 along the Main Central Thrust (e.g., Srivastava and Mitra, 1994; Hodges, 2000;
1071 Mitra et al., 2010; Tobgay et al., 2012; Law et al., 2013) and quartz OH contents
1072 (50 - 150 ppm) far lower than required for water weakening challenge our
1073 understanding of dislocation creep and the role of water in deformation deep in the
1074 continental crust. With only few exceptions, IR spectra of Main Central Thrust
1075 samples have OH bands of the same character and size as dry natural quartz
1076 crystals, which are strong and have not been deformed by dislocation processes in
1077 laboratory experiments (e.g., Heard and Carter, 1968; Blacic, 1975; Blacic and
1078 Christie, 1984). Given sufficiently high temperatures, it is possible that quartz
1079 may deform at tectonic strain rates without critical hydrogen defects at
1080 dislocations and water weakening (Kilian et al., 2016). However, this implies that
1081 we have not measured flow laws for appropriately dry quartzites that we can apply
1082 to amphibolite conditions and natural strain rates. Alternatively, water may have
1083 been lost from quartz interiors following deformation. It is also possible that
1084 hydrogen defects that enhance dislocation motion at high temperatures and natural
1085 strain rates may be sourced from grain boundaries or micas, diffusing over longer
1086 distances than are possible at greenschist conditions or laboratory strain rates.

1087 Spatial variations in OH content of quartz in natural shear zones, as
1088 mapped in this study, may provide key insights into the role of water weakening,
1089 and changes in water content during deformation and recrystallization. High
1090 resolution FTIR imaging of OH in Moine Thrust and Main Central Thrust samples
1091 show that water contents are increased, fluid inclusions are redistributed, and
1092 water contents are decreased during brittle deformation, plastic creep, recovery,
1093 and recrystallization. Changes in OH contents of quartz mylonites and the history
1094 of fluid migration during deformation may lead to changes in governing flow
1095 laws, non-steady rates of creep, and shifting zones of localized shear.

1096

1097 **6 Conclusions**



1098

1099 The brightness of synchrotron IR radiation enables measurement of IR spectra for
1100 much smaller sampling volumes than is possible using conventional global IR
1101 sources of FTIR instruments. In this study, we have used this improvement in
1102 signal to characterize and measure small OH absorption bands in quartz mylonites
1103 with an aperture size of 10 μm and to map water contents spatially. The ability to
1104 measure IR spectra for small, individual recrystallized grains by methods
1105 described in this study is limited by samples that must be $> 40 \mu\text{m}$ thick to avoid
1106 internal reflections that lead to interference fringes that are larger than OH
1107 absorption bands. High-resolution images of OH in quartz mylonites, based on
1108 spectra collected through a 10 μm aperture as samples are translated under the
1109 beam, reveal large variations in OH content that correspond to the distributions of
1110 fluid inclusions and layer silicates, and to deformation and recrystallization
1111 microstructures. The OH contents of quartz in Moine Thrust mylonites deformed
1112 at greenschist conditions are comparable to wet quartzites deformed in the
1113 laboratory by processes of water weakening. By comparison, OH contents of
1114 quartz in Main Central Thrust mylonites deformed at upper amphibolite conditions
1115 are small, and molecular water, as required to deform quartz at experimental strain
1116 rates, is absent. What role water plays in deformation at these conditions is
1117 unclear, calling for further studies of water weakening in natural shear zones.
1118 High-resolution FTIR mapping of OH offers a new method of tracking changes in
1119 water content during deformation, recovery, and recrystallization.

1120

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1140

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1661

1662 **Table Captions**

1663

1664 **Table 1.** Structural distances of samples below the Moine Thrust at the Stack of
1665 Glencoul (Law et al., 1986, 2010) and above the Main Central Thrust on the NW
1666 and Eastern Suttlej transects (Law et al., 2013). Deformation temperatures
1667 estimated by Law et al. (2013) using the Kruhl (1998) quartz *c*-axis fabric opening
1668 angle thermometer, and temperatures and pressures of metamorphism estimated by
1669 Stahr (2013) using THERMOCALC multi-equilibria thermometry (Powell and
1670 Holland 1994) are indicated.

1671

1672 **Figure Captions**

1673

1674 **Figure 1.** Doubly polished IR plates of Moine Thrust mylonites prepared
1675 perpendicular to foliation and parallel to lineation. (a) Low magnification image
1676 of large IR plate prepared from Stack of Glencoul sample SG-10 (unpolarized
1677 light imaged in high resolution Nikon slide scanner), with a mean thickness $t =$
1678 120 μm . (b) Same IR plate of SG-10 shown in (a) but with crossed-polarized
1679 light. The Moine Thrust transport direction is shown in the plane of the IR plate



1680 (top to the left, in the geographic WNW direction). (c) Higher magnification
1681 optical micrograph of large IR plate of sample SG-10, with crossed polarized light
1682 and local plate thickness (117 μm) determined from IR interference fringes. (d)
1683 Optical micrograph of Stack of Glencoul sample SG-7 in crossed-polarized light
1684 (normal 30 μm section thickness), illustrating deformation and recover
1685 microstructures with higher resolution than in thick IR plates, with undulatory
1686 extinction in ribbon quartz grains, fine recrystallized grains and aligned muscovite
1687 grains localized at quartz grain boundaries. (e) Low magnification image of
1688 ultrathin IR plate of Stack of Glencoul sample SG-7-5 (crossed-polarized light)
1689 mounted on a copper TEM slot ring ($t = 4\text{--}8 \mu\text{m}$, based on IR interference fringes).
1690 (f) Higher magnification optical micrograph of the same IR plate SG-7-5 (crossed-
1691 polarized light) as shown in (e) with deformation and recrystallization
1692 microstructures shown more clearly than in normal 30 μm thin section.

1693

1694 **Figure 2.** IR spectra of quartz grains in mylonites from the Moine Thrust (Stack
1695 of Glencoul) and the Main Central Thrust (Sutlej Valley) using conventional FTIR
1696 with globar IR source through a 100 μm aperture. (a) Absorbance spectra of
1697 quartz grains of Moine Thrust IR plate SG-10 (shown in Fig. 1, $t = 120 \mu\text{m}$)
1698 showing broad OH band at 3400 cm^{-1} associated with fluid inclusions in milky
1699 quartz. (b) Absorption spectra of Moine Thrust quartz grains of the same IR plate
1700 (SG-10) showing the same broad OH band due to fluid inclusions in addition to an
1701 OH absorption band (or shoulder) at 3600 cm^{-1} due to mica inclusions. (c)
1702 Absorbance spectra of quartz grains of Main Central Thrust IR plates S09-30 ($t =$
1703 $121 \mu\text{m}$), S09-35 ($t = 74 \mu\text{m}$), and S09-58 ($t = 145 \mu\text{m}$), with many grains showing
1704 small sharp OH absorption bands at 3595 , 3482 , 3431 , 3408 , and 3379 cm^{-1} due to
1705 hydrogen point defects, and less common grains with a large broad absorption
1706 band at 3400 cm^{-1} . (d) Absorption spectra of quartz grains of Main Central Thrust
1707 sample S09-30 ($t = 121 \mu\text{m}$) with a prominent OH band at 3600 cm^{-1} due to mica
1708 inclusions and smaller OH bands at lower wavenumbers. Absorbance values



1709 plotted vertically are normalized to represent their values for a uniform sample
1710 thickness of 1 mm.
1711
1712 **Figure 3.** OH absorption bands due to fluid inclusions and hydrogen defects in
1713 quartz grains and those due to mica inclusions. (a) Integrated absorbances of OH
1714 bands at 3400 cm^{-1} (due to fluid inclusions) and 3595 to 3379 cm^{-1} (due to
1715 hydrogen interstitials) of quartz grains ($\Delta_{\text{OH}_{\text{qtz}}}$ in cm^{-2}) can be separated from the
1716 integrated absorbance of the 3600 cm^{-1} OH band of mica inclusions ($\Delta_{\text{OH}_{\text{mica}}}$),
1717 depending on the IR absorption spectra of the micas. The integrated absorbance Δ
1718 of the 3600 cm^{-1} OH band of micas can be determined over 3702 cm^{-1} to 3544 cm^{-1} ,
1719 fitting a straight line between these values. In order to evaluate the remaining
1720 OH absorbance of the spectrum, the value of Δ determined for the 3600 cm^{-1} band
1721 is deducted from value from the integrated absorbance Δ due to all OH absorptions
1722 over 3705 cm^{-1} to 2880 cm^{-1} , again fitting the background by a straight line. This
1723 procedure works well to distinguish OH absorbances due to muscovite inclusions
1724 ($\Delta_{\text{OH}_{\text{musc}}}$) from those of fluid inclusions and hydrogen defects of Moine Thrust
1725 quartz grains ($\Delta_{\text{OH}_{\text{qtz}}}$ determined for SG-10, Grain 2, $t = 120\text{ }\mu\text{m}$), given that (b)
1726 the absorption spectrum of muscovite grains in Moine Thrust samples is very
1727 simple, with one prominent OH absorption band at 3620 cm^{-1} ($t = 8\text{ }\mu\text{m}$,
1728 absorbance normalized to $t = 0.1\text{ mm}$). OH absorption spectra of (c) muscovite
1729 and (d) biotite grains of Main Central Thrust mylonites (S09-63, $t = 155\text{ }\mu\text{m}$) are
1730 more complex than displayed by muscovite grains of the Moine Thrust. The
1731 prominent OH band of muscovite grains in Main Central Thrust samples appears
1732 at 3638 cm^{-1} with smaller OH bands at 3815 , 3311 , 3146 , and 3035 cm^{-1} . The
1733 primary OH band of biotite grains in Main Central Thrust samples appears at 3614
1734 cm^{-1} (with shoulders at 3674 and 3561 cm^{-1}) and secondary OH bands at 3258 ,
1735 3043 , and 2829 cm^{-1} . The OH bands of all micas at $\sim 3600\text{ cm}^{-1}$ are highly
1736 anisotropic; absorbances at $\sim 3600\text{ cm}^{-1}$ for muscovite grains measured in
1737 polarized IR radiation are maximized when the vibration direction E is parallel to



1738 the basal (001) plane, while principal OH absorbances for biotite grains are
1739 maximized when E is perpendicular to (001). However, secondary OH bands of
1740 muscovite and biotite grains of the Main Central Thrust show significant
1741 variations from the dichroism displayed by the primary OH bands of each of these
1742 micas. The complex nature of these secondary OH bands presents significant
1743 difficulties in separating OH absorptions of quartz grains due to molecular water
1744 and hydrogen defects and those due to mica inclusions dispersed within quartz
1745 grains. All absorption spectra shown here were measured using unpolarized IR
1746 radiation with a conventional FTIR spectrometer and a 100 μm aperture.
1747 Absorbances for all spectra but (b) are normalized to a uniform thickness t of 1
1748 mm.

1749

1750 **Figure 4.** IR spectra of quartz grains in Moine Thrust mylonite samples (Stack of
1751 Glencoul), measured with synchrotron-FTIR system using a 10 μm aperture with
1752 varying IR plate thicknesses (a) $t = 113 \mu\text{m}$ (BSG 10, local plate thickness
1753 determined from interference fringes), (b) $t = 13 \mu\text{m}$ (SG-10.2t), and (c) $t = 6.5$
1754 μm (BSG 7.3, sample plate SG-7-1). For doubly polished sample plates of ~ 100
1755 μm , IR spectra measured using the coupled synchrotron-FTIR spectrometer
1756 system through a 10 μm aperture are similar to those measured using a
1757 conventional FTIR spectrometer through a 100 μm aperture. (a) IR spectra of
1758 Moine Thrust sample (SG-10) show OH absorption bands of similar character at
1759 the same wavenumbers for a 10 μm aperture as OH bands measured through a
1760 larger (100 μm) aperture, including a large broad absorption band at 3400 cm^{-1} due
1761 to dispersed fluid inclusions (both BSG 10.1 and BSG 10.3) and a sharper band at
1762 3600 cm^{-1} due to mica inclusions (shown by BSG 10.3). Interference fringes in
1763 samples $\sim 100 \mu\text{m}$ thick are apparent, allowing determination of local IR plate
1764 thickness, but they do not obscure the OH absorption bands. (b) Interference
1765 fringes for samples $< 25 \mu\text{m}$ thick are large, and make detection of small OH
1766 absorption bands difficult. The only detectable absorbance bands in sample SG-



1767 10.2t ($t = 13 \mu\text{m}$) are due to strong primary Si-O vibrations (at $\nu < 2200 \text{ cm}^{-1}$);
1768 OH absorption bands cannot be resolved. (c) Interference fringes are very large
1769 for thin IR plates ($t = 6.5 \mu\text{m}$; SG-7-1); OH absorption bands are not observed,
1770 even after attempts to model them and remove fringes numerically. All
1771 absorbance values (and apparent absorbance values of interference fringes
1772 exhibited by SG-7) are normalized to a uniform sample thickness of 1 mm.
1773

1774 **Figure 5.** Ultrathin IR plates ($t = 4\text{--}8 \mu\text{m}$) exhibit deformation microstructures in
1775 crossed-polarized light with greater clarity than those of conventional thin sections
1776 ($t = 30 \mu\text{m}$). (a) Deformed ribbon quartz grains of Moine Thrust sample SG-7-6 (t
1777 $= 6\text{--}8 \mu\text{m}$) show distinct subgrains with sharply defined subgrain walls, while
1778 extinction in conventional thin sections of ribbon quartz is smoothly varying and
1779 subgrain boundaries are difficult to detect. Finely recrystallized quartz grains at
1780 sutured grain boundaries are smaller than subgrains within large deformed quartz
1781 grains. (b) Less deformed quartz porphyroblast of Moine Thrust sample SG-7-5 (t
1782 $= 5\text{--}6 \mu\text{m}$) with internal subgrains, and surrounding recrystallized quartz grains of
1783 similar dimensions to those of internal subgrains.
1784

1785 **Figure 6.** OH absorbance maps of Moine Thrust sample SG-7 constructed from
1786 900 IR spectra (SG-7t-1map, 30×30 steps, $10 \mu\text{m}/\text{step}$) for a doubly polished
1787 plate prepared ($56 \mu\text{m}$ thick) perpendicular to foliation and parallel to the transport
1788 direction (lineation horizontal in all panels). (a) Plane-polarized micrograph of IR
1789 plate SG-7t showing nearly equant globular augen quartz grain surrounded by
1790 highly deformed ribbon and recrystallized quartz grains. Light scattering is mostly
1791 due to micas and fluid inclusions. Outlined box is the region imaged by integrated
1792 IR absorbances. (b) The same plane-polarized light micrograph as in (a) with
1793 superposed map of OH absorbance of the 3600 cm^{-1} band due to micas (OH of
1794 micas given in integrated area cm^{-2}). Contours are given in $\log_{10}(\Delta \text{ in } \text{cm}^{-2})$ for the
1795 integrated absorbance of the $\sim 3600 \text{ cm}^{-1}$ OH band. (c) Cross-polarized light



1796 micrograph of the same region of IR plate SG-7t as shown in (a), with subtle
1797 undulatory extinction of ribbon quartz shown in white and first-order interference
1798 colors. (d) The same cross-polarized light micrograph as in (c) with superposed
1799 map of OH absorbance of the broad 3400 cm^{-1} band due to molecular water in
1800 fluid inclusions (and smaller OH bands due to hydrogen defects) of quartz grains
1801 (OH of quartz in molar ppm, $\text{OH}/10^6\text{ Si}$). Contours given in \log_{10} (ppm) for
1802 integrated absorbance of the broad 3400 cm^{-1} OH band (and lesser sharp OH
1803 bands) of quartz. OH contents of quartz and those associated with micas are
1804 contoured in color with blue (and cool colors) corresponding to large water (OH)
1805 contents, and red (and warm colors) corresponding to low water (OH) contents.
1806

1807 **Figure 7.** OH absorbance maps of Moine Thrust sample SG-10 constructed from
1808 3600 IR spectra (SG-10-XY2, 60×60 steps, $10\text{ }\mu\text{m}/\text{step}$) for a doubly polished
1809 plate prepared ($117\text{ }\mu\text{m}$ thick) perpendicular to foliation and parallel to the
1810 transport direction (lineation horizontal in all panels). (a) Plane-polarized
1811 micrograph of IR plate SG-10-2 showing deformed ribbon quartz grains
1812 surrounded by recrystallized quartz grains. Light scattering is mostly due to micas
1813 and fluid inclusions. Outlined box is the region imaged by integrated IR
1814 absorbances. (b) The same plane-polarized light micrograph as in (a) with
1815 superposed map of OH absorbance of the 3600 cm^{-1} band due to micas (OH of
1816 micas given in integrated area $\Delta\text{ cm}^{-2}$ and contoured as $\log_{10}\Delta$). (c) Cross-
1817 polarized light micrograph of the same region of IR plate SG-10-2 as shown in (a),
1818 with undulatory extinction of deformed ribbon quartz grains, incipient
1819 recrystallized grains at quartz ribbon margins, and recrystallized matrix grains
1820 shown by first-order interference colors. (d) The same cross-polarized light
1821 micrograph as in (c) with superposed map of OH absorbance of the broad 3400
1822 cm^{-1} band due to molecular water in fluid inclusions (and smaller OH bands due to
1823 hydrogen defects) of quartz grains (OH of quartz in molar ppm, $\text{OH}/10^6\text{ Si}$ and
1824 contoured as $\log_{10}\text{ OH}$). OH contents of quartz and those associated with micas



1825 are contoured in color with blue (and cool colors) corresponding to large water
1826 (OH) contents, and red (and warm colors) corresponding to low water (OH)
1827 contents.

1828

1829 **Figure 8.** OH absorbance maps of Moine Thrust sample SG-8 constructed from
1830 1800 IR spectra (SG-8t-map1, 30 x 60 steps, 10 $\mu\text{m}/\text{step}$) for a doubly polished
1831 plate prepared (73 μm thick) perpendicular to foliation and parallel to the transport
1832 direction (lineation horizontal in all panels). (a) Plane-polarized micrograph of IR
1833 plate SG-8t showing deformed ribbon quartz grains and regions of finely dispersed
1834 mica and recrystallized quartz grains. Light scattering is mostly due to micas,
1835 fluid inclusions, and grain boundaries. Outlined box is the region imaged by
1836 integrated IR absorbances. (b) The same plane-polarized light micrograph as in
1837 (a) with superposed map of OH absorbance of the 3600 cm^{-1} band due to micas
1838 (OH of micas given in integrated area Δcm^{-2} and contoured as $\log_{10} \Delta$). (c) Cross-
1839 polarized light micrograph of the same region of IR plate SG-8t as shown in (a),
1840 with undulatory extinction of deformed ribbon quartz grains and regions of
1841 recrystallized grains shown by first-order interference colors. (d) The same cross-
1842 polarized light micrograph as in (c) with superposed map of OH absorbance of the
1843 broad 3400 cm^{-1} band due to molecular water in fluid inclusions (and smaller OH
1844 bands due to hydrogen defects) of quartz grains (OH of quartz in molar ppm,
1845 $\text{OH}/10^6 \text{ Si}$ and contoured as $\log_{10} \text{OH}$). OH contents of quartz and those
1846 associated with micas are contoured in color with blue (and cool colors)
1847 corresponding to large water (OH) contents, and red (and warm colors)
1848 corresponding to low water (OH) contents.

1849

1850 **Figure 9.** OH absorbance maps of Main Central Thrust sample S09-63
1851 constructed from 1200 IR spectra (S09-63-map1, 30 x 40 steps, 50 $\mu\text{m}/\text{step}$) for a
1852 doubly polished plate prepared (155 μm thick) perpendicular to foliation and
1853 parallel to the transport direction (lineation horizontal in all panels). (a) Plane-



1854 polarized micrograph of IR plate S-09-63 showing coarse, clear quartz grains, and
1855 coarse muscovite and biotite grains with readily distinguishable color and
1856 pleochroism. Scattering of light is primarily due to grain boundaries, with little
1857 evidence for dense fluid inclusions or finely dispersed micas. Outlined box is the
1858 region imaged by integrated IR absorbances. (b) The same plane-polarized light
1859 micrograph as in (a) with superposed map of OH absorbance of the 3600 cm^{-1}
1860 band due to micas (OH of micas given in integrated area $\Delta\text{ cm}^{-2}$ and contoured as
1861 $\log_{10} \Delta$). (c) Cross-polarized light micrograph of the same region of IR plate SG-
1862 09-63 as shown in (a), with high-temperature deformation and recovery
1863 microstructures (in higher order interference colors) that are characterized by
1864 subtle (to absent) undulatory extinction of quartz, subgrain walls, and coarse
1865 recrystallized grains. (d) The same cross-polarized light micrograph as in (c) with
1866 superposed map of OH absorbance over 3705 and 2880 cm^{-1} to capture the broad
1867 and sharp bands of quartz, deducting the large 3600 cm^{-1} band of micas, but
1868 including smaller OH bands of micas between 3311 and 2920 cm^{-1} . Contours in
1869 this absorbance can only be attributed unambiguously to fluid inclusions and
1870 hydrogen defects in quartz where micas (and their 3600 cm^{-1} absorbances) are
1871 absent. OH contents of quartz and those associated with micas are contoured in
1872 color with blue (and cool colors) corresponding to large water (OH) contents, and
1873 red (and warm colors) corresponding to low water (OH) contents.



Table 1 Synchrotron FTIR Imaging of OH Kronenberg et al

Moine Thrust - Stack of Glencoul mylonitic quartzites

Sample	Distance below MT
SG.7	2.5 m
SG.8	2.9 m
SG.10	4.6 m

Main Central Thrust - NW Suttlej transect orthogneisses

Sample	Distance above MCT	T
S09-30	~ 750 m	~ 600°C ¹
S09-35	75 m	~ 540°C ¹

Main Central Thrust - Eastern Suttlej transect paragneisses

Sample	Distance above MCT	T	P
S09-58	~ 4,500 m	735°C ²	900 MPa ²
S09-63	~1500 m	675°C ²	850 MPa ²
S09-71B	25 m	610°C ¹	

¹Deformation temperature estimated from quartz *c*-axis fabric opening angle (Law et al. 2013).

²Temperature and pressure of metamorphism estimated from THERMOCALC multi-equilibra thermobarometry (Stahr 2013, p. 67; Law et al., 2013).



Figures

Synchrotron FTIR Imaging of OH

Kronenberg et al

Fig. 1a

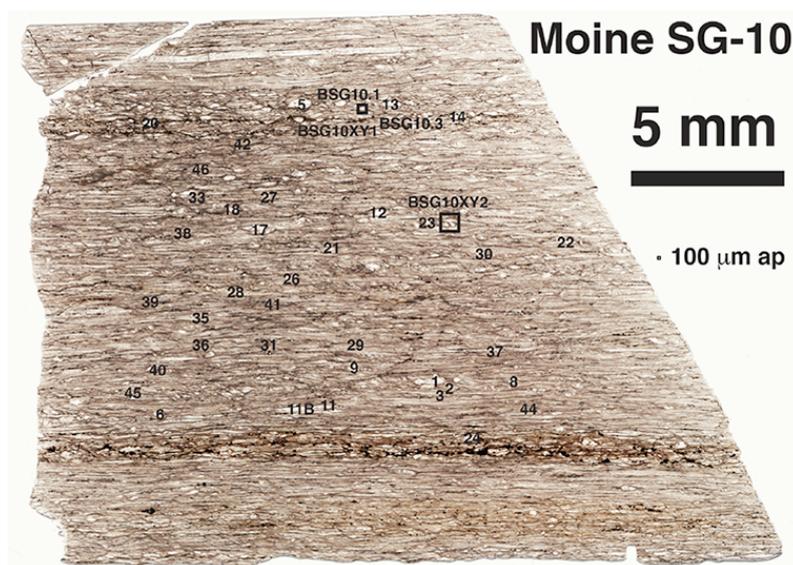


Fig. 1b

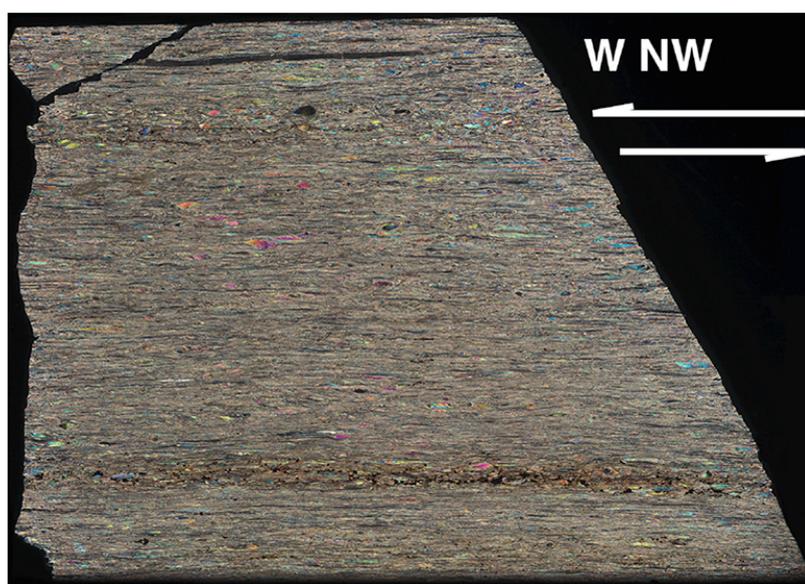




Fig. 1c

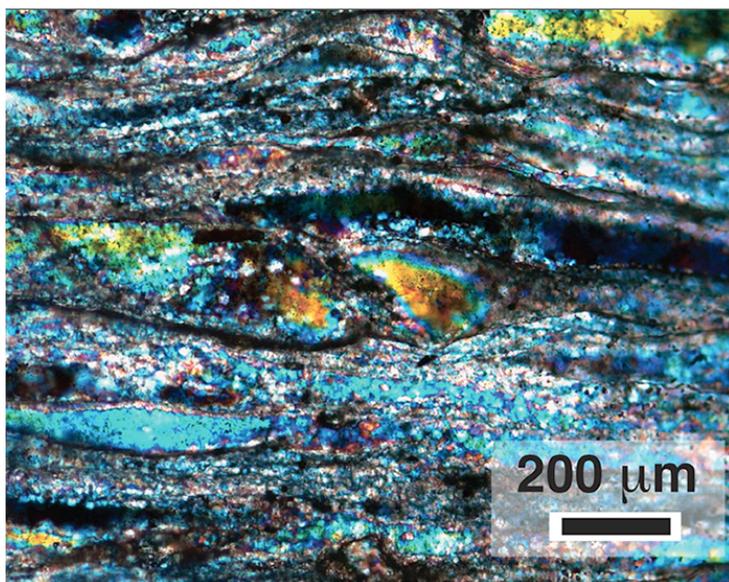


Fig. 1d

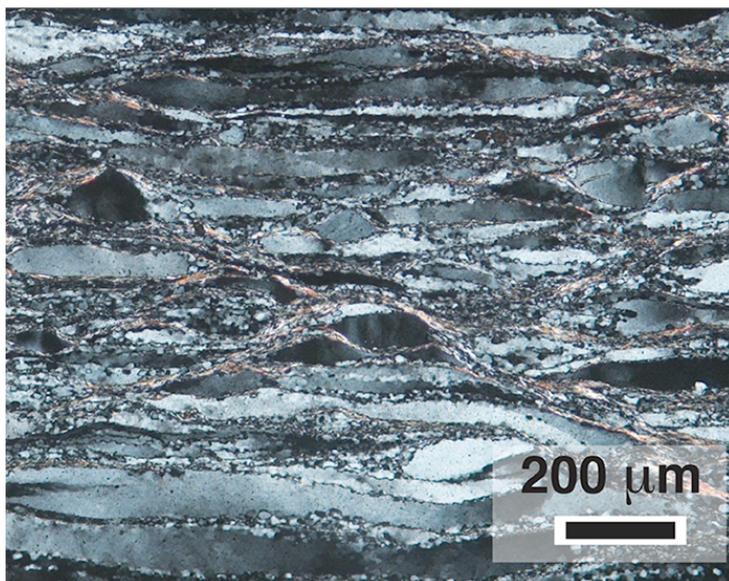




Fig. 1e

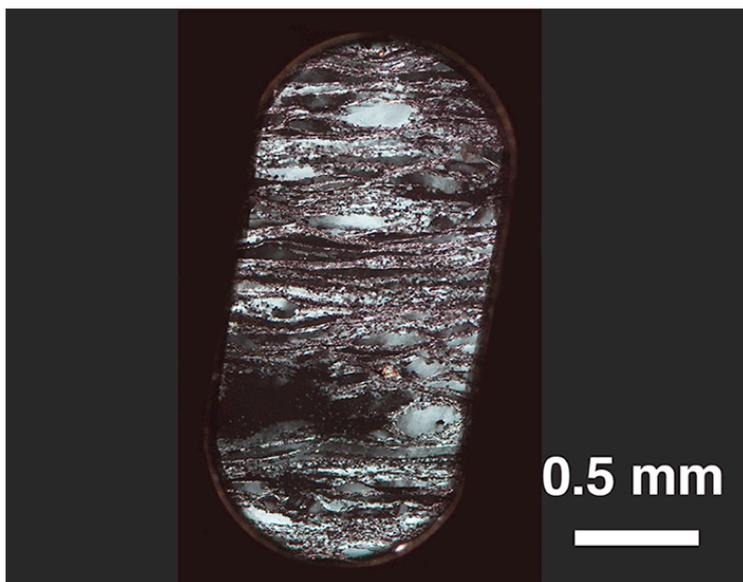


Fig. 1f

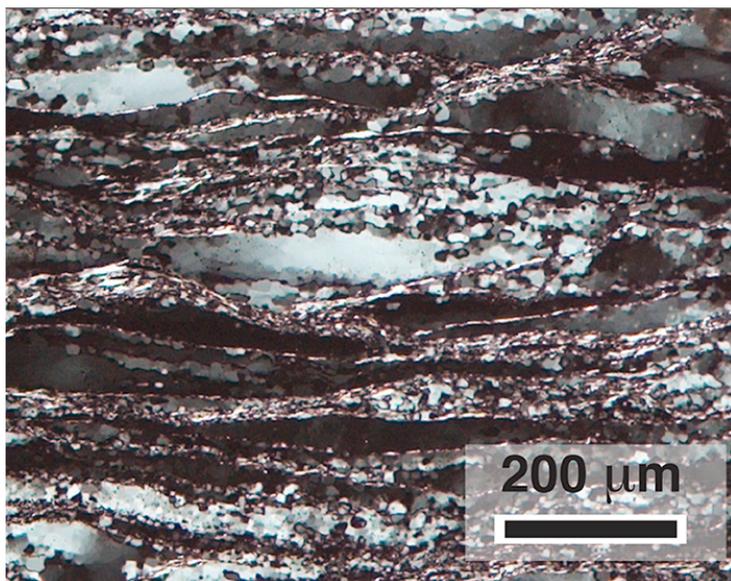




Fig. 2a

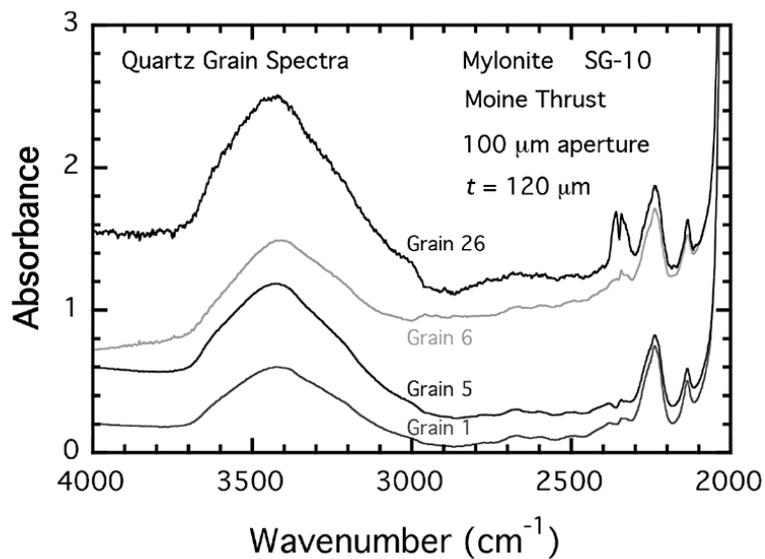


Fig. 2b

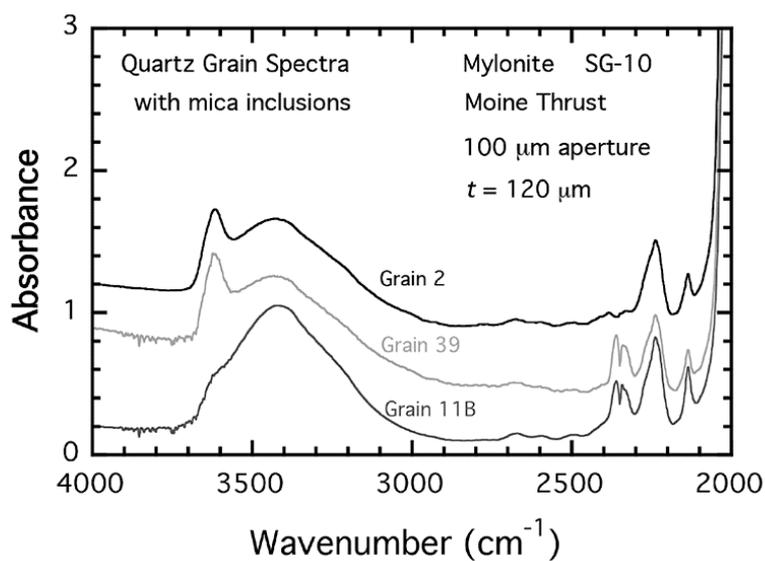




Fig. 2c

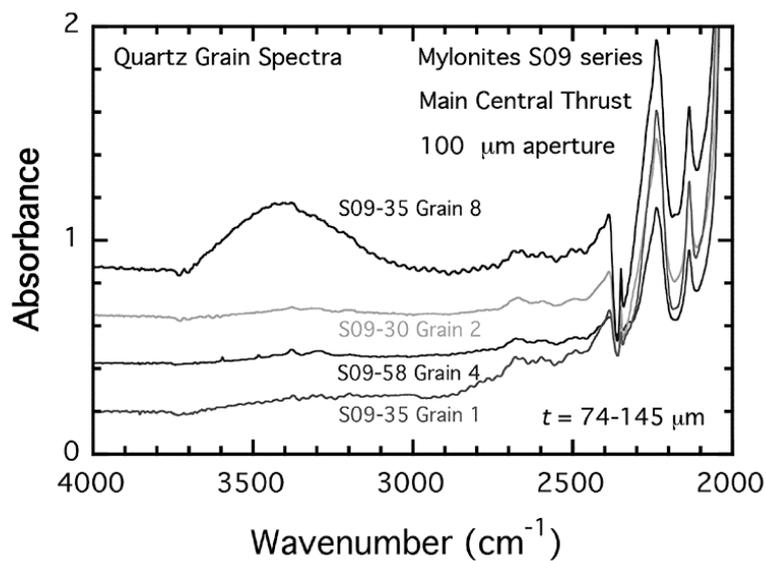


Fig. 2d

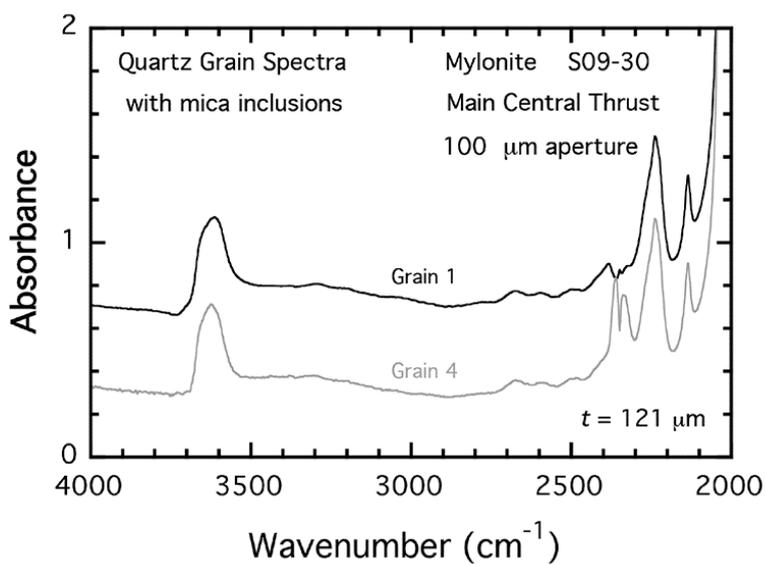




Fig. 3a

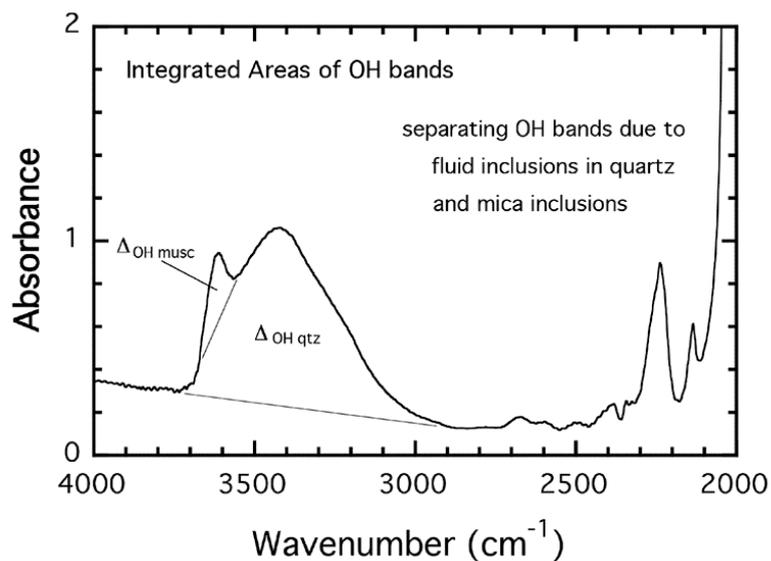


Fig. 3b

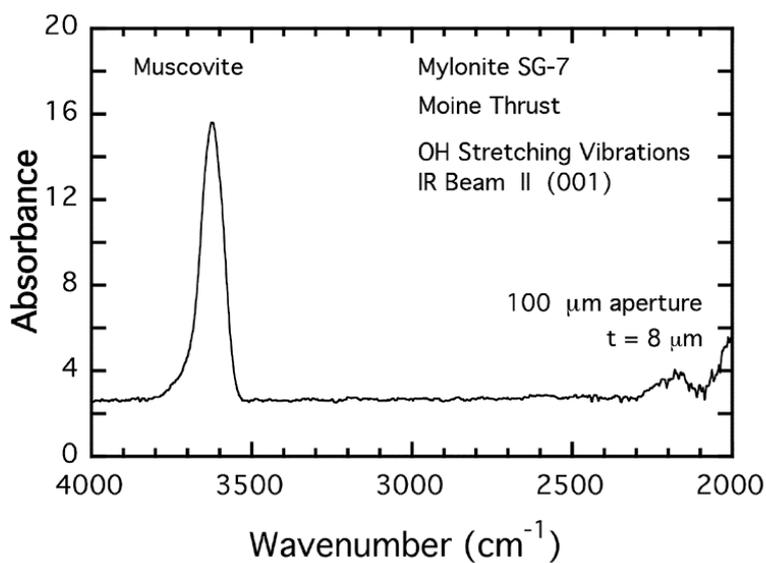




Fig. 3c

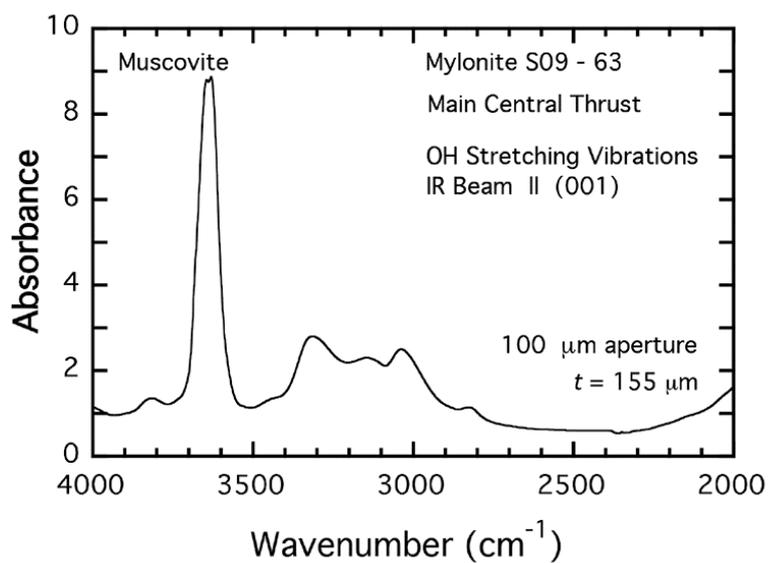


Fig. 3d

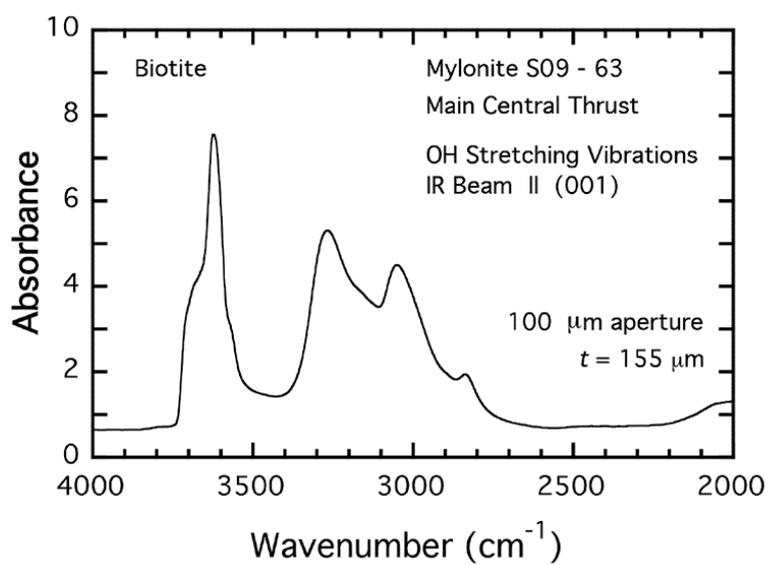




Fig. 4a

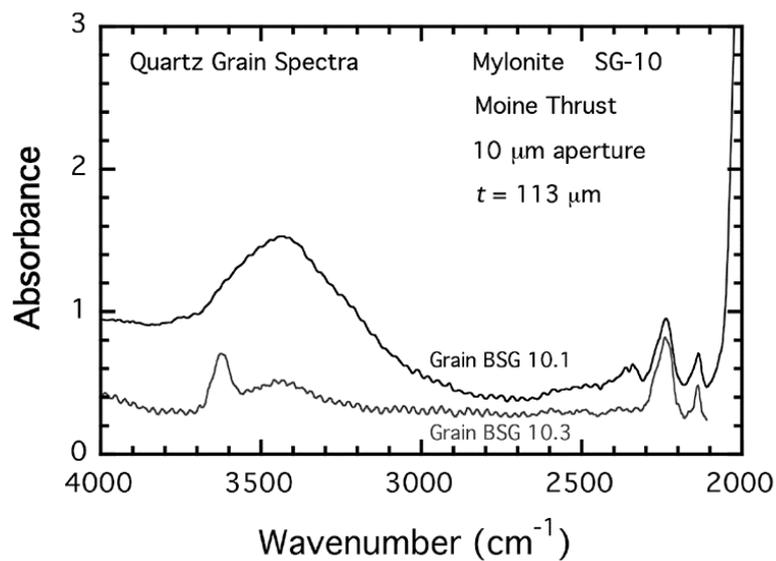


Fig. 4b

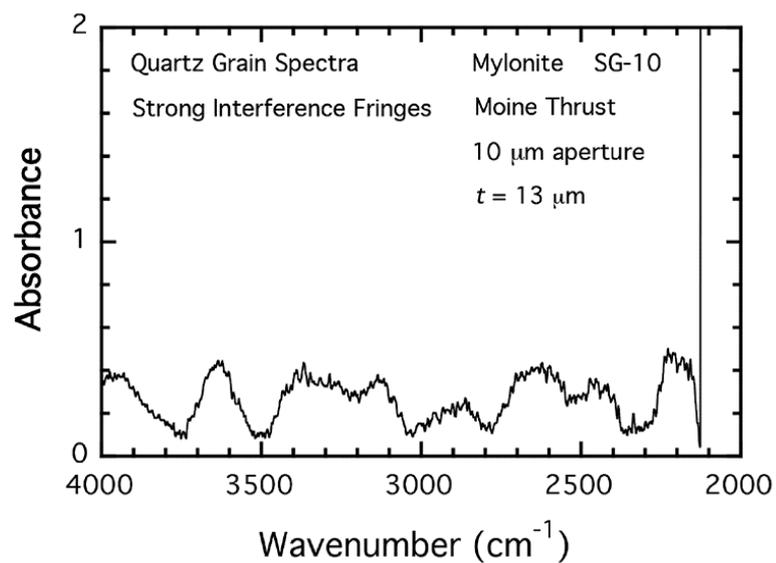




Fig. 4c

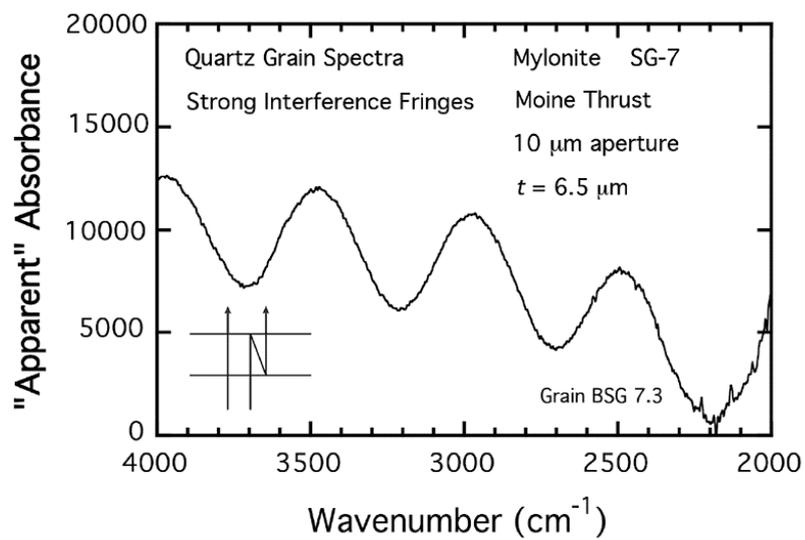




Fig. 5a

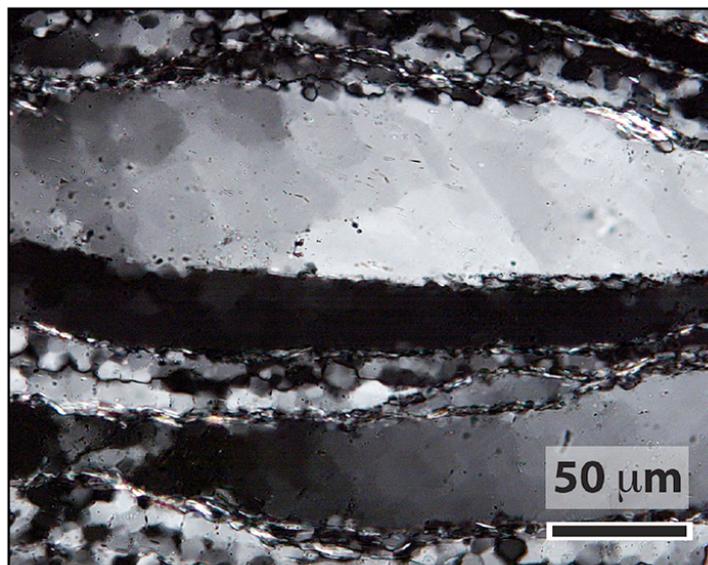


Fig. 5b

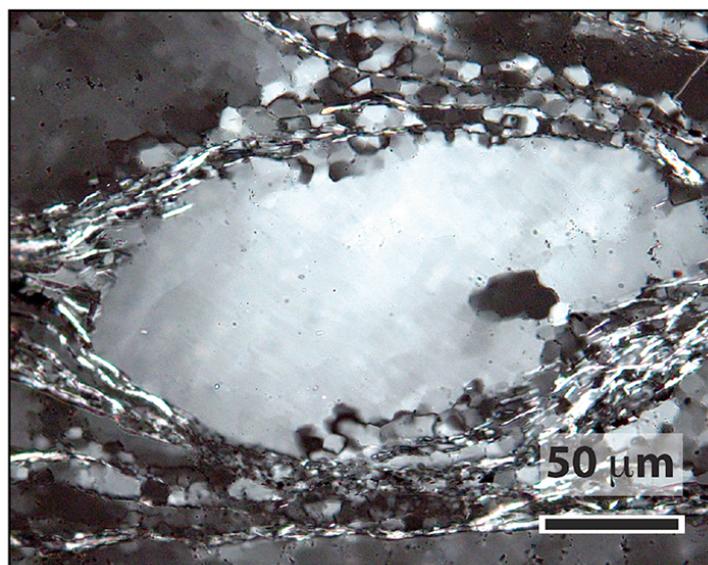
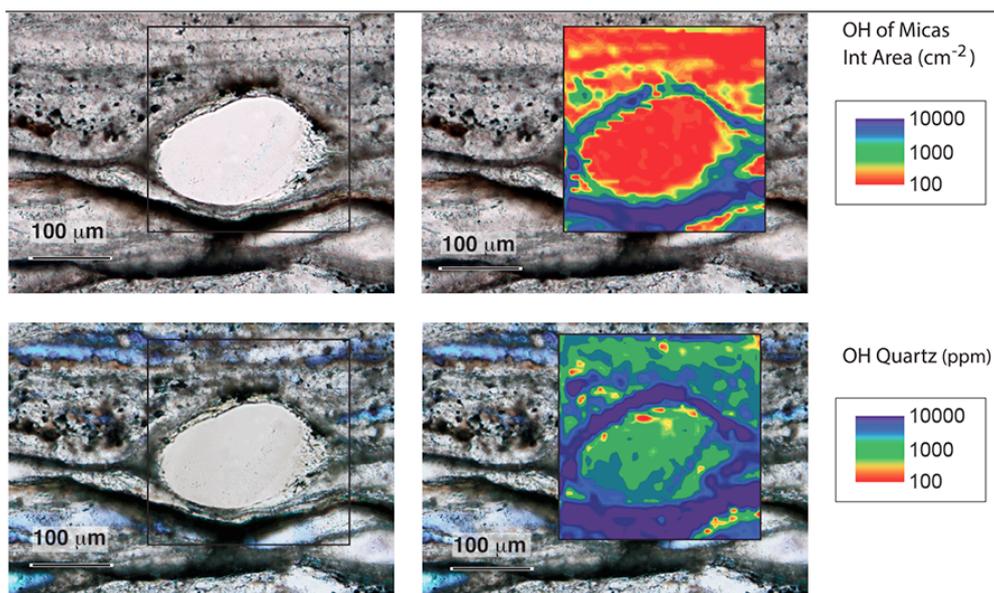


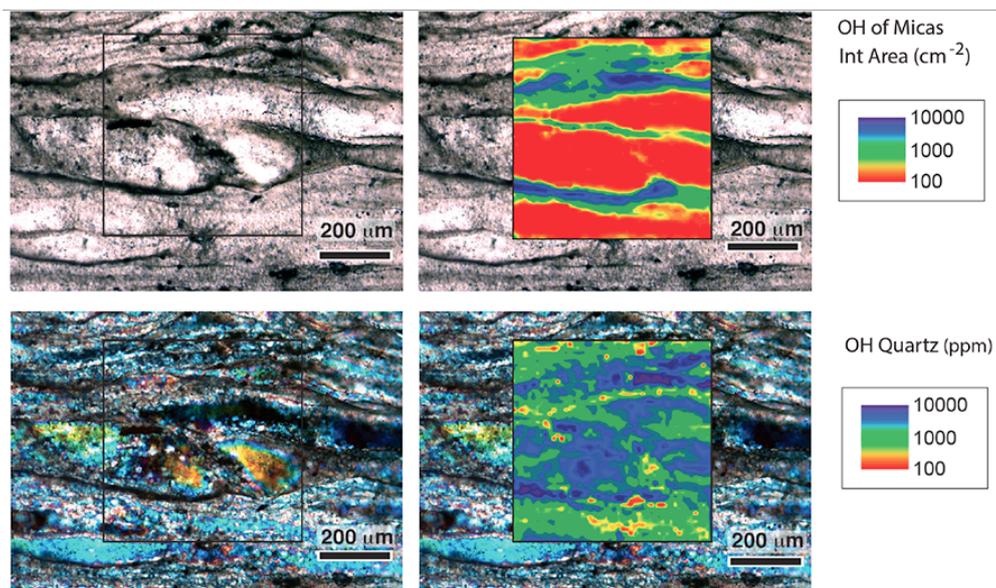


Fig. 6



SG-7

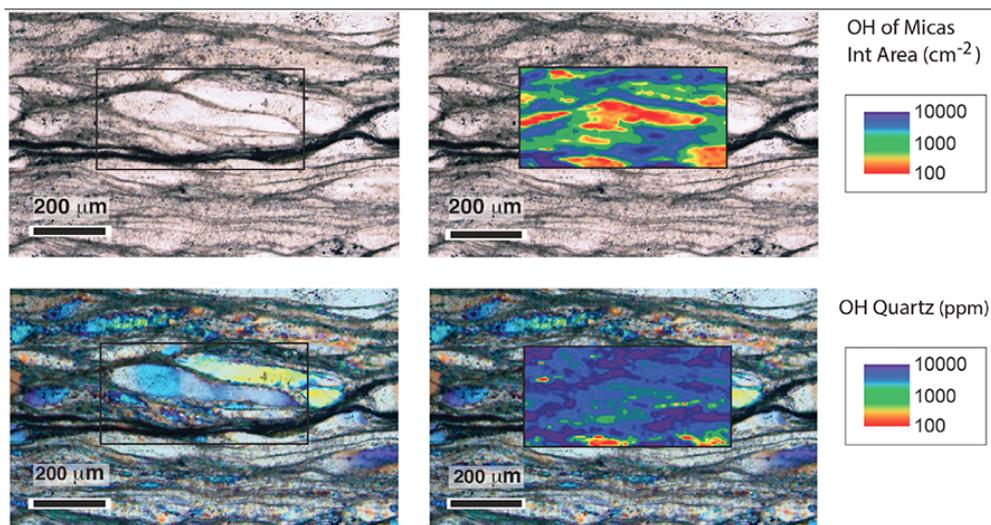
Fig 7



SG-10

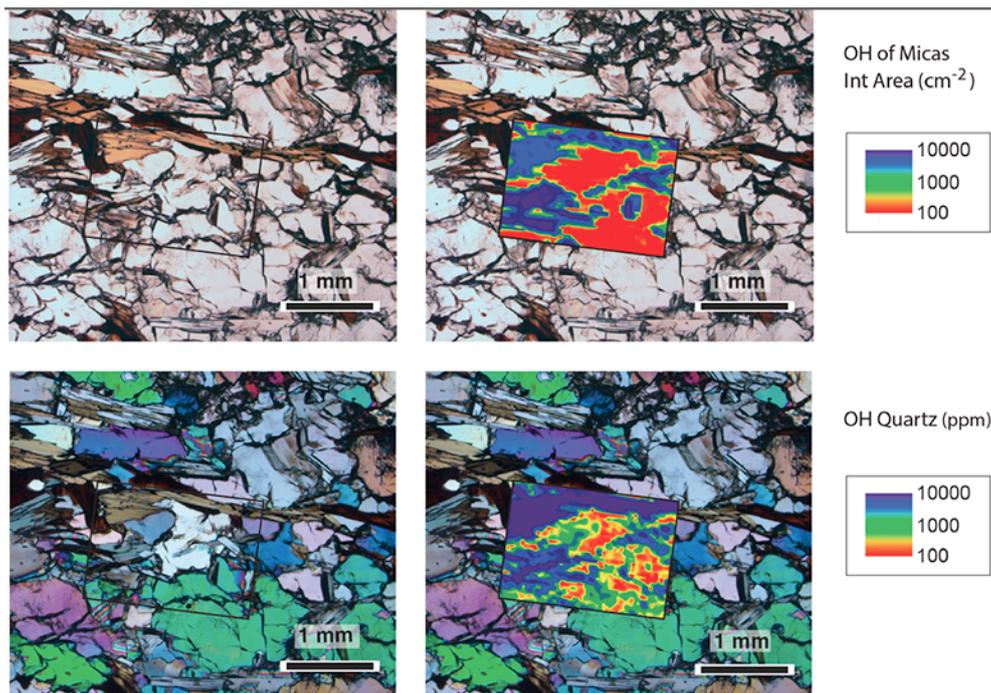


Fig. 8



SG-8

Fig. 9



S09-63